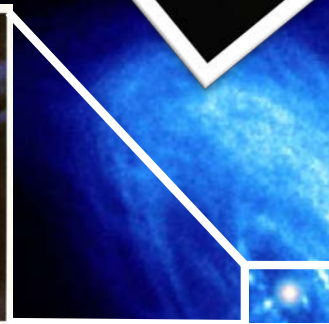
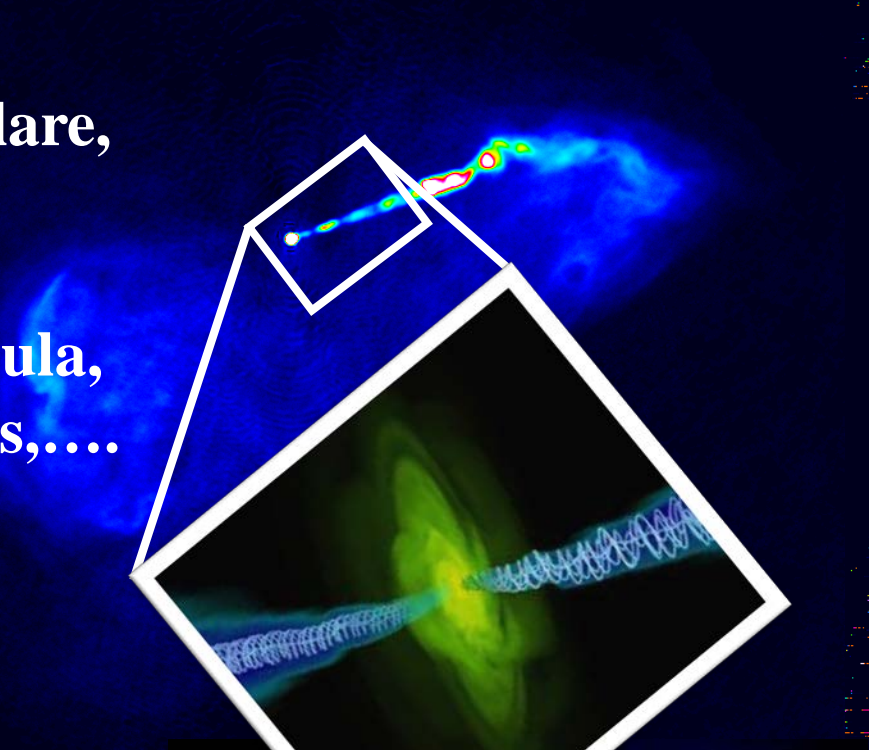


Thermodynamical Approach of Energy Partition During Magnetic Reconnection

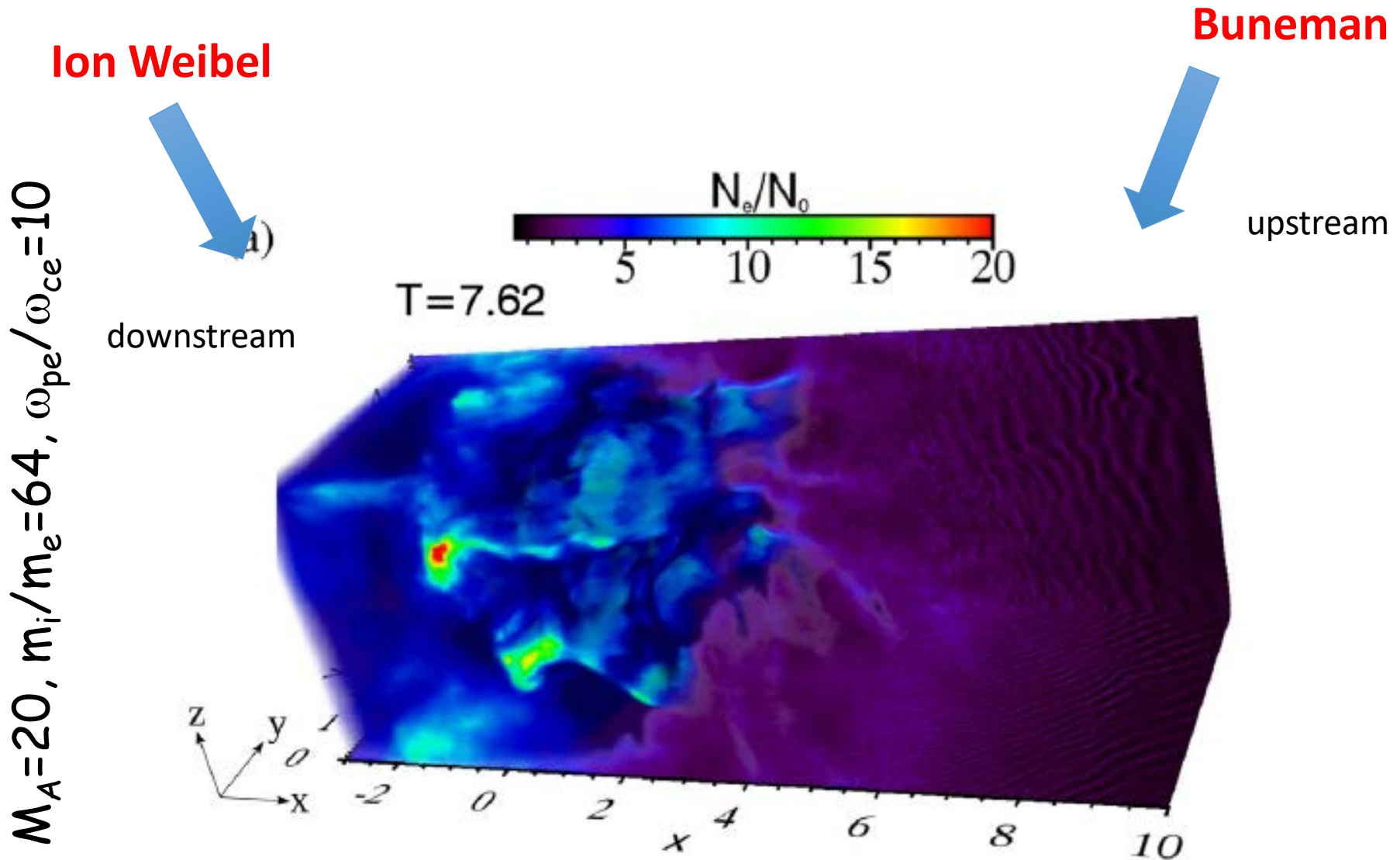
Masahiro Hoshino
University of Tokyo

**Solar & Stellar Flare,
Magnetosphere,
Accretion Disks,
Pulsar Wind-Nebula,
Astrophysical Jets,....**



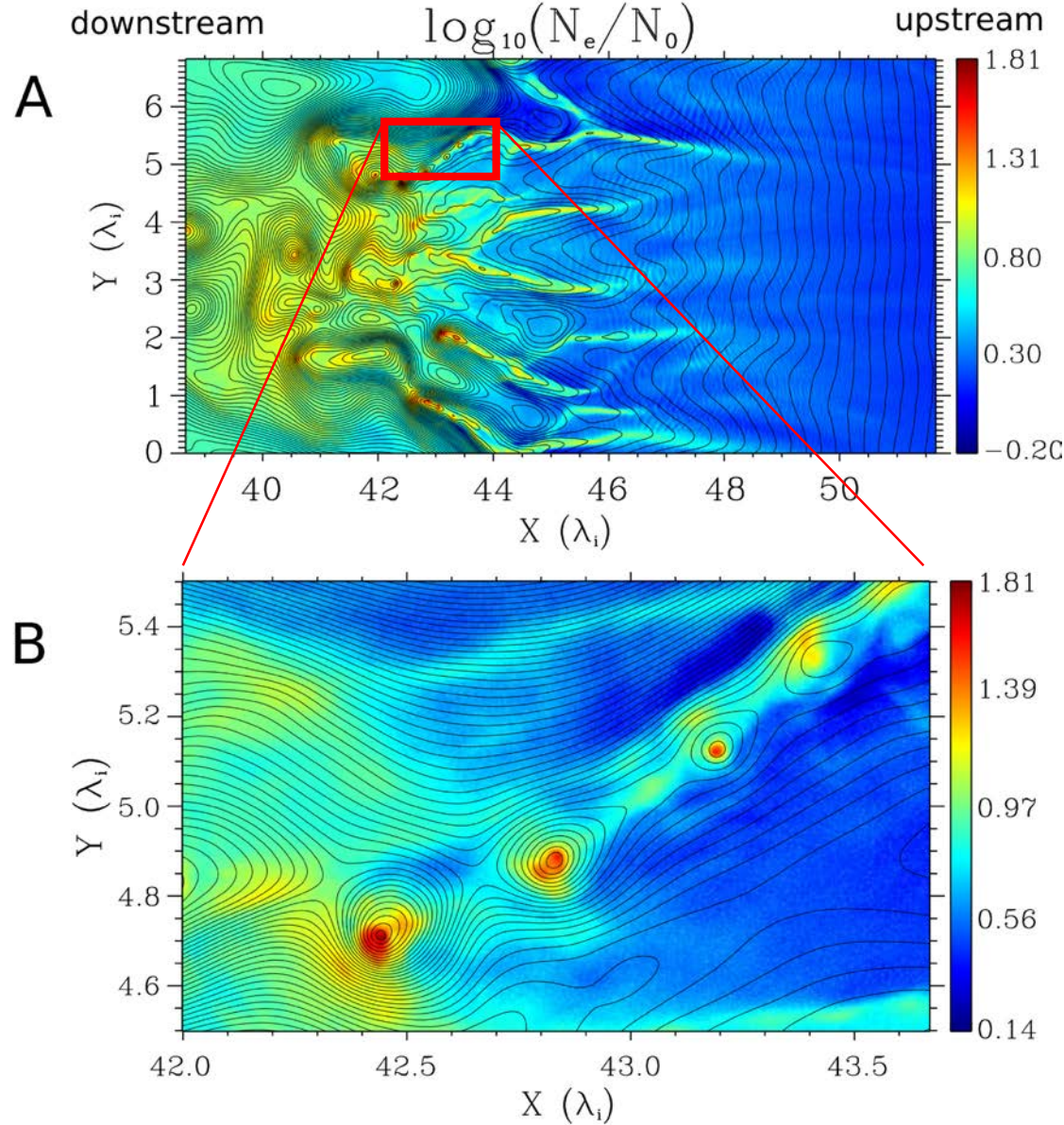
**Rapid Energy Dissipation &
Nonthermal Particle Acceleration**

3D High Mach Number Shock



Ion Weibel and Magnetic Reconnection

$$M_A=42, m_i/m_e=225, \omega_{pe}/\omega_{ce}=10$$



fast flow



Ion reflection from shock front



Temperature anisotropy

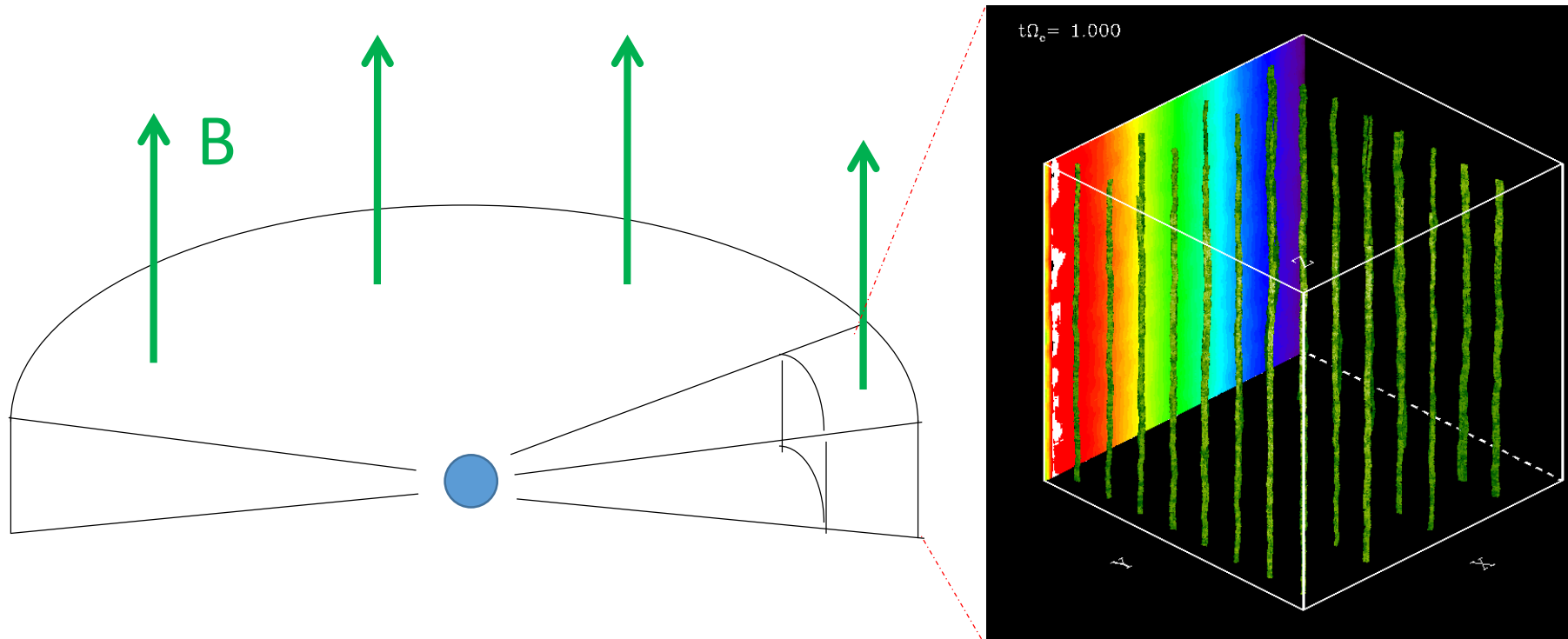


Weibel instability & B field generation



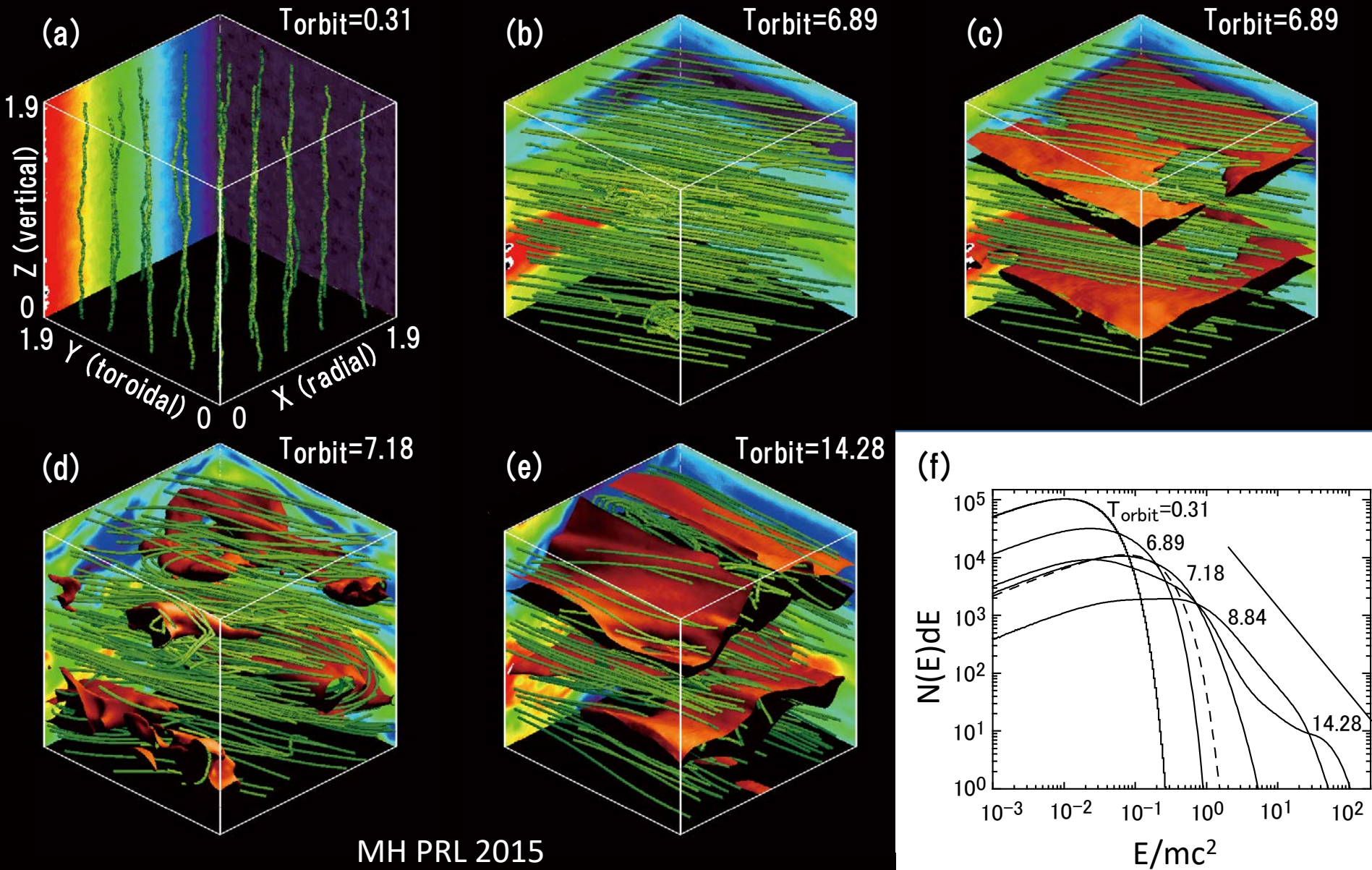
Magnetic reconnection

MRI and Reconnection in PIC simulation



$\beta=1536$, Kepler rotation Ω
300³ grids 40 particles/cell,
periodic shearing box, electron-positron plasma

Particle Acceleration in Accretion Disks

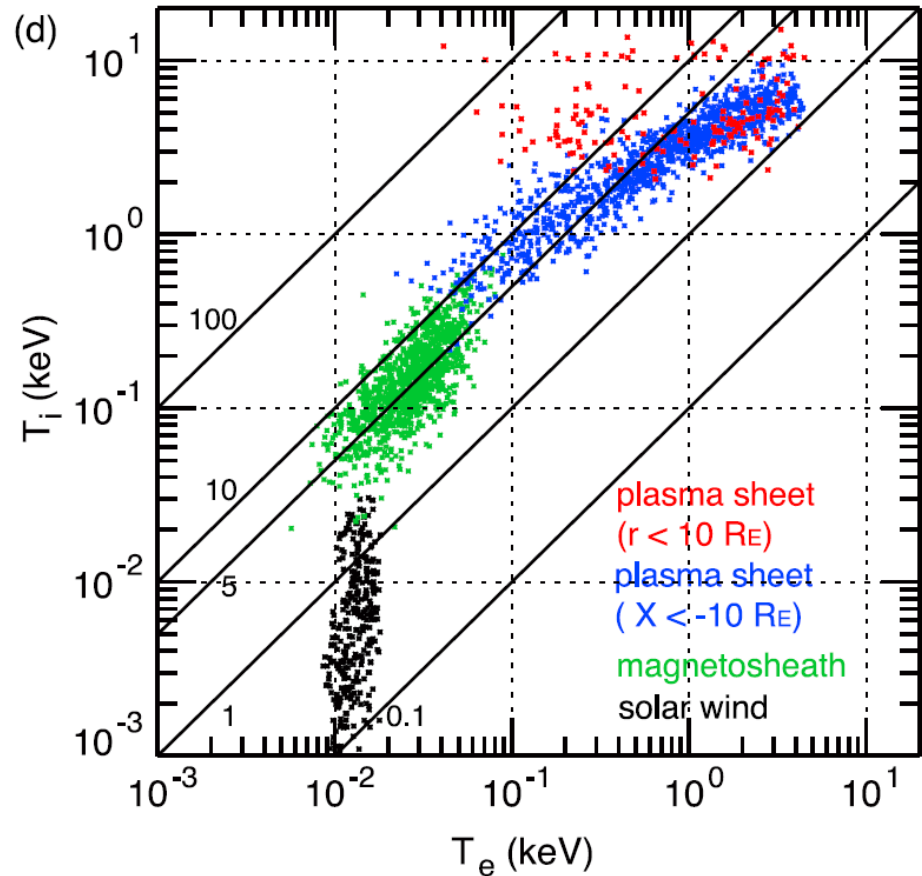


Observations of T_i/T_e

magnetosphere

$$T_i/T_e = 5 \sim 10$$

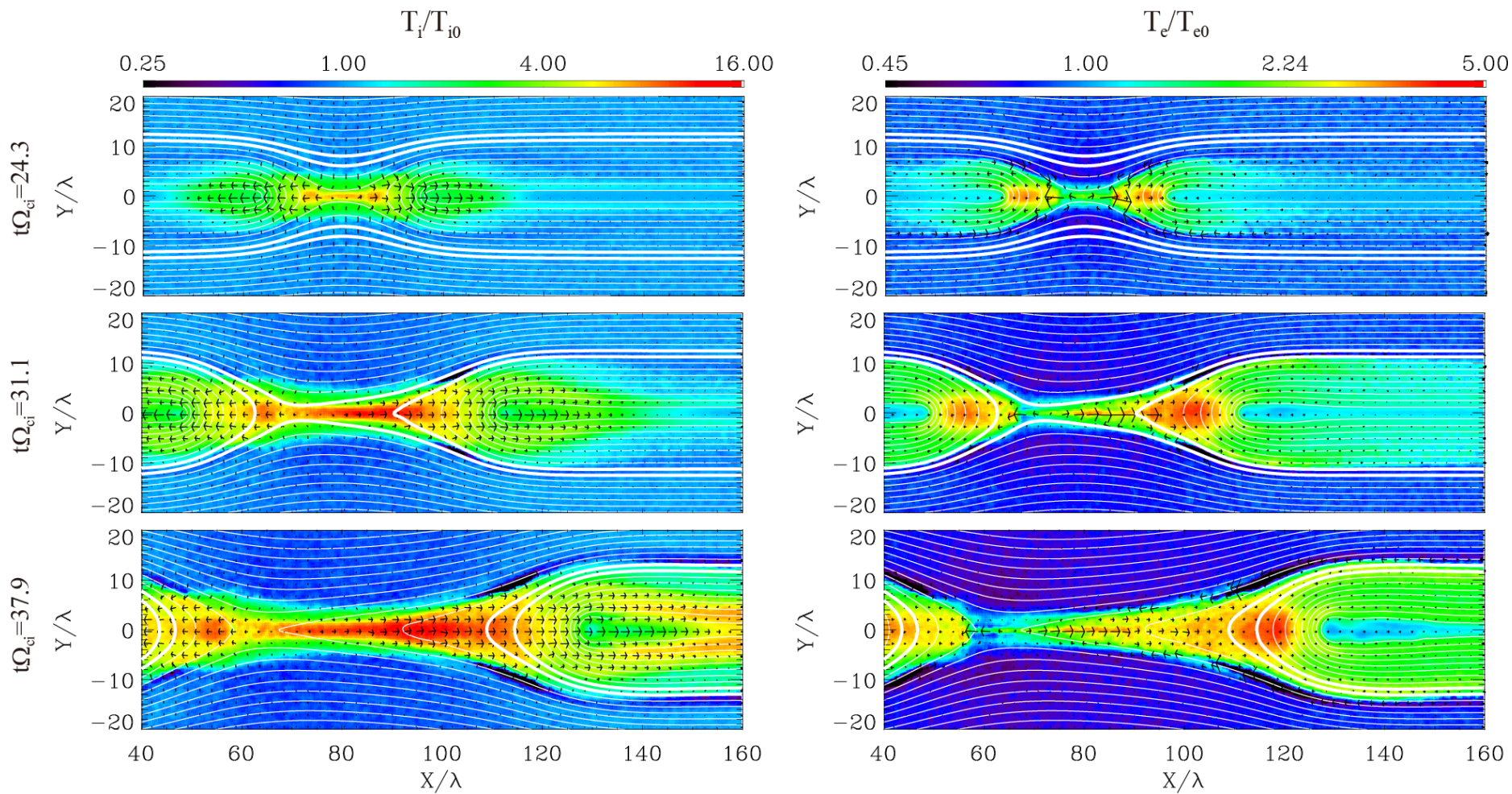
Hot ions are believed to be generated during magnetic reconnection...



(cf. Baumjohann+ JGR 1989; Eastwood+ PRL 2013; Phan+ GRL 2013)

Wang+ JGR 2012

T_i & T_e Heating in PIC simulation



Motion of flux tube in 2D

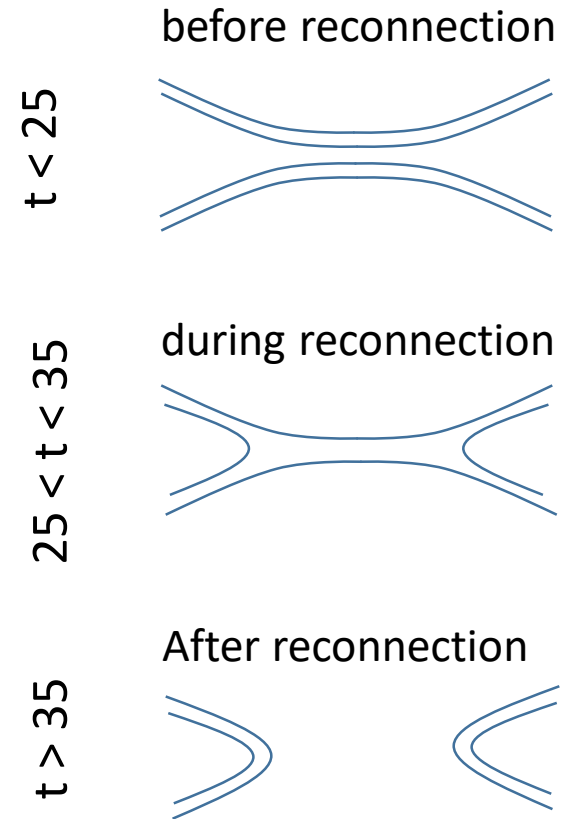
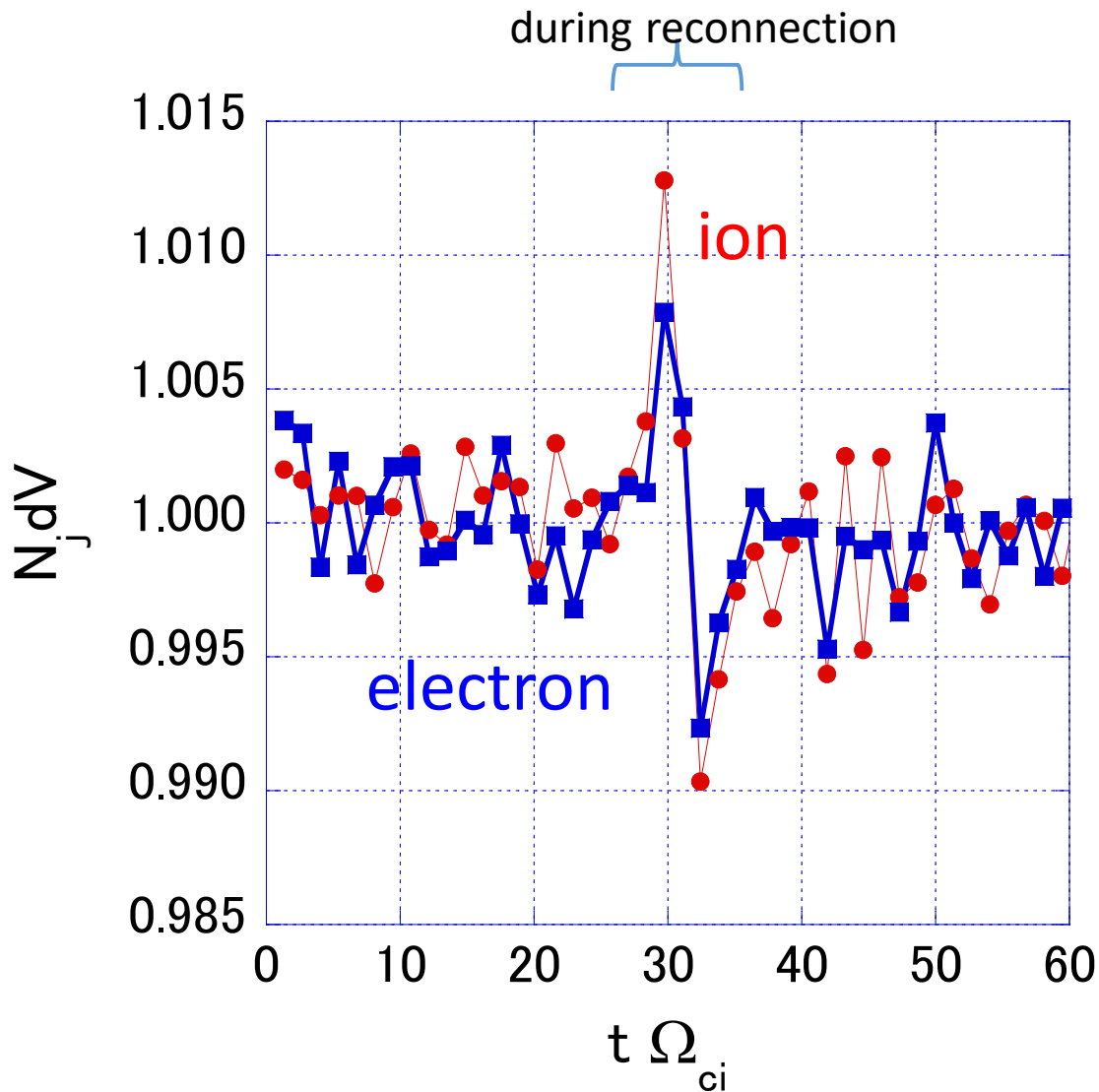
$$\vec{B}(x, y) = \nabla \times A_z(x, y)\vec{e}_z + B_z(x, y)\vec{e}_z$$

$$\frac{dx}{B_x(x, y)} = \frac{dy}{B_y(x, y)} \iff dA_z(x, y) = 0$$

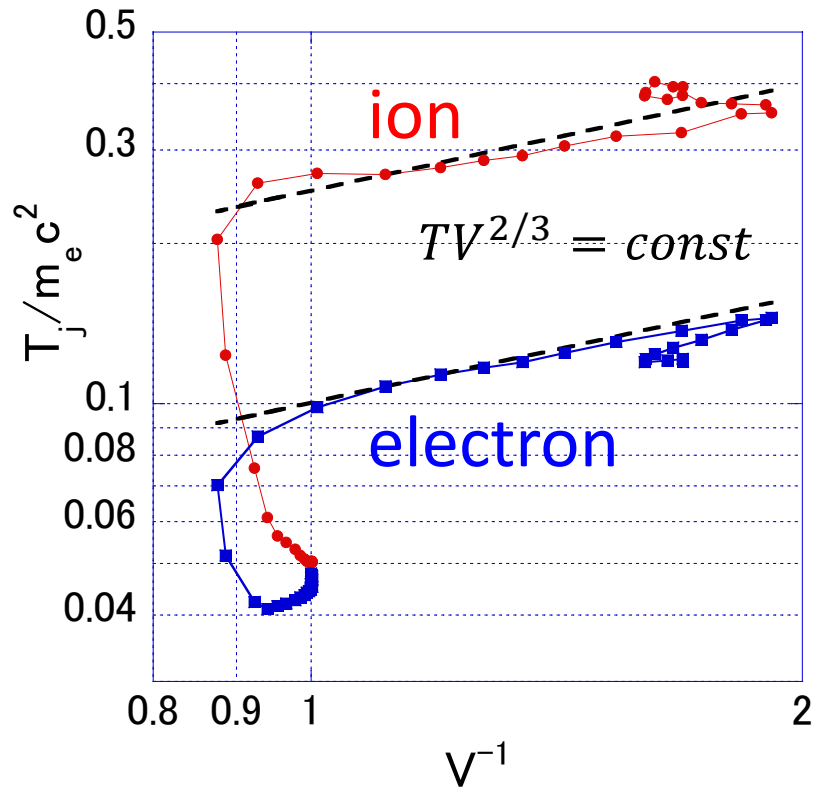
$$\text{If } \vec{E} + \frac{1}{c}\vec{v} \times \vec{B} = 0,$$

$$\text{then } \left(\frac{\partial}{\partial t} + \vec{v} \cdot \nabla \right) A_z(x, y, t) = 0.$$

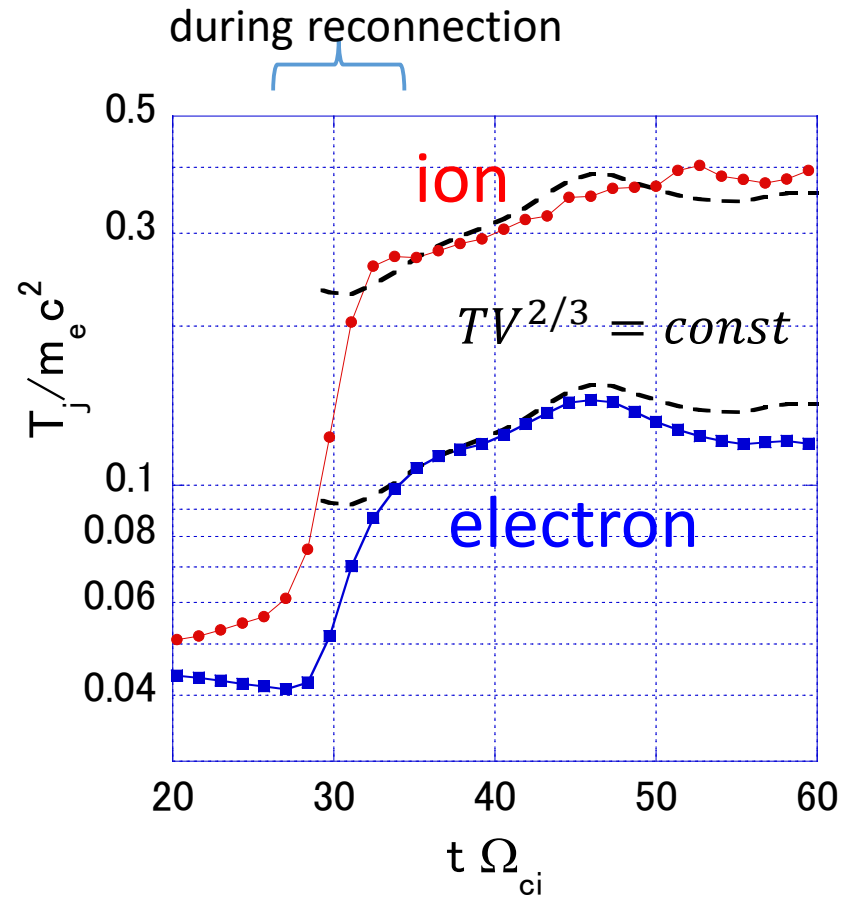
Time History of N in Flux Tube



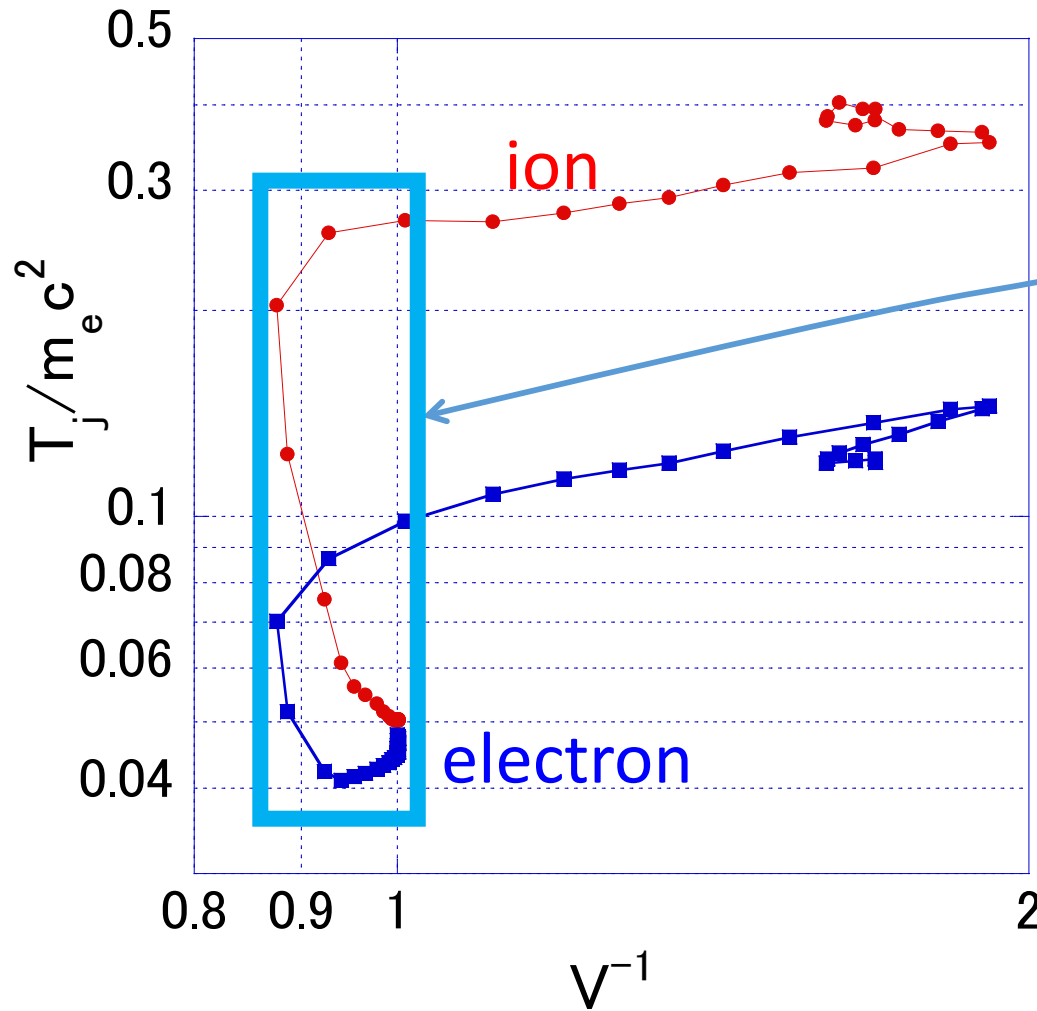
T-V Relations



(V : Volume of Flux Tube)

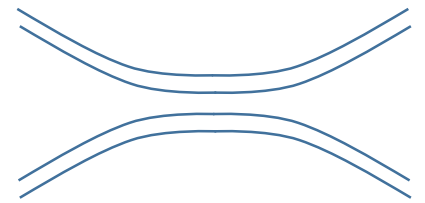


T - V relation

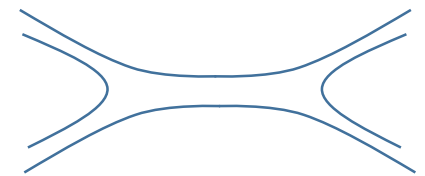


(V : Volume of Flux Tube)

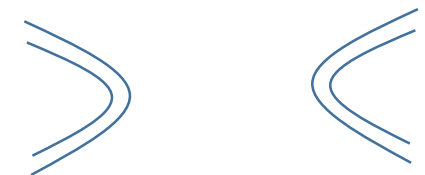
before reconnection



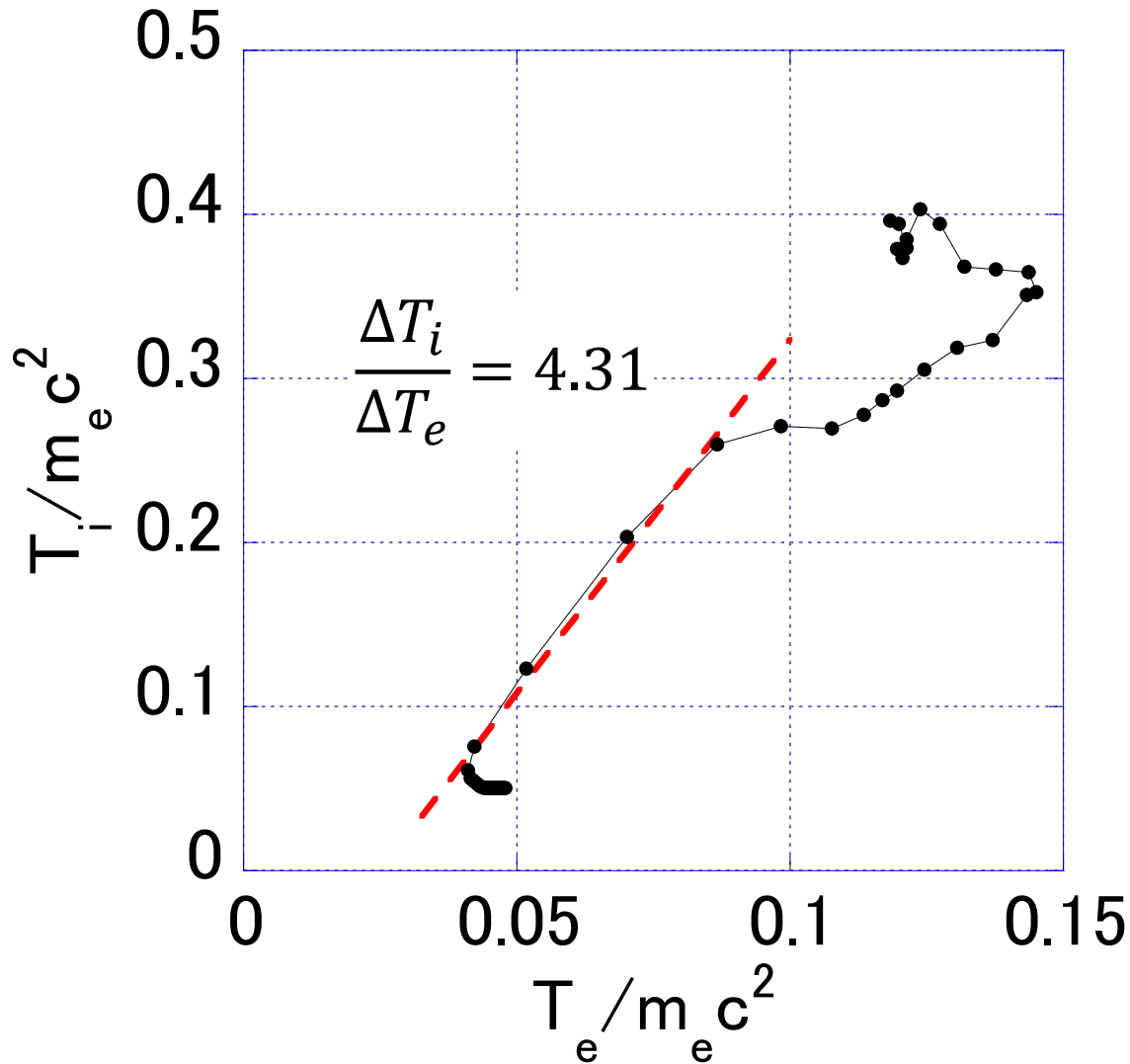
during reconnection



After reconnection



Time history of T_i and T_e

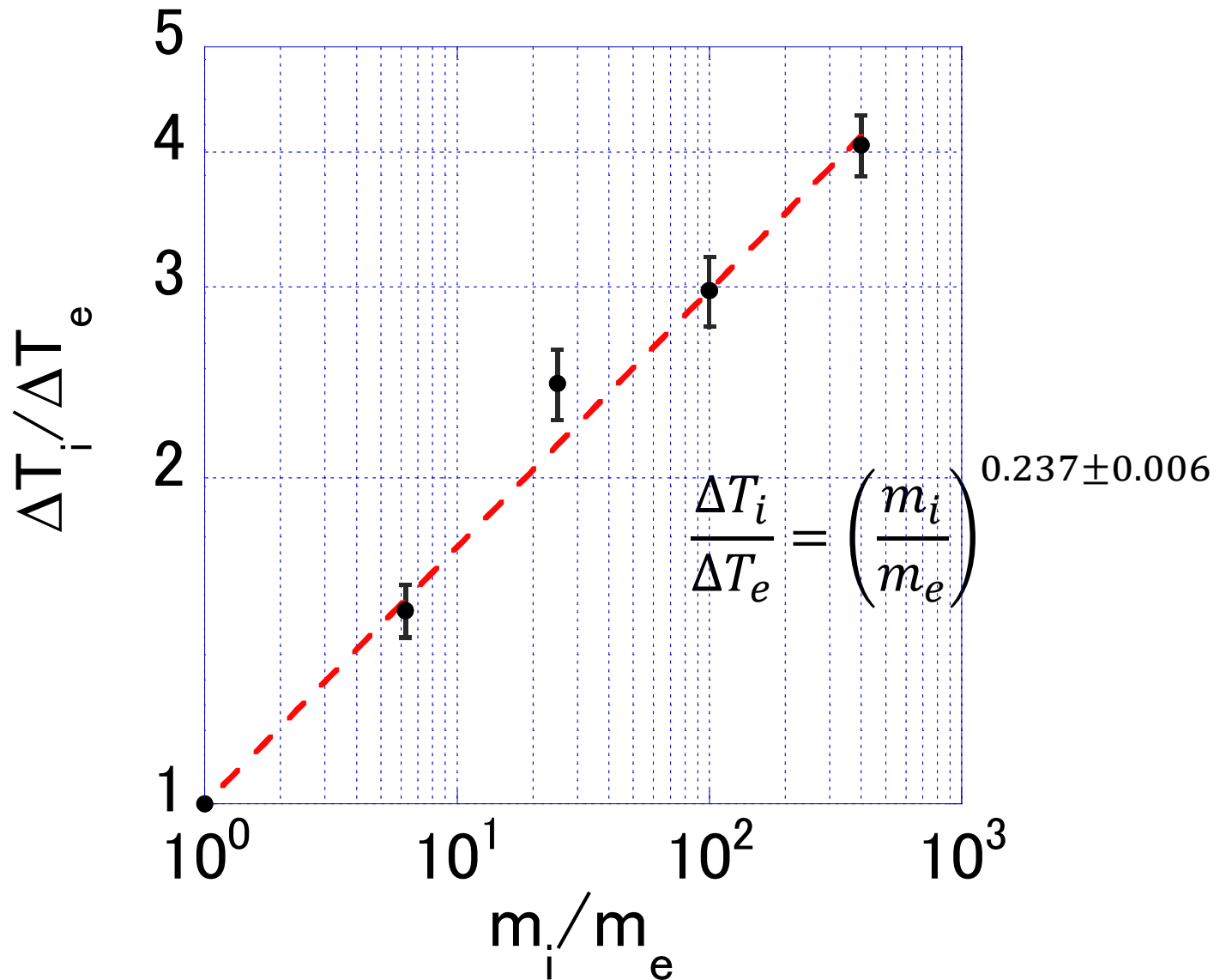


$$\frac{m_i}{m_e} = 400$$

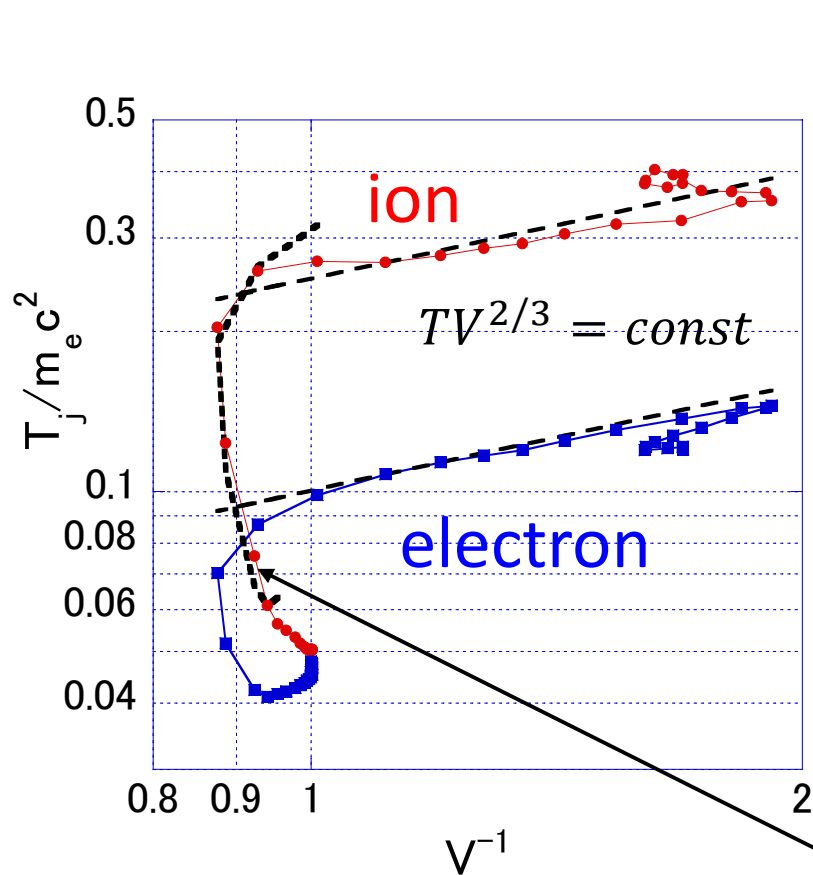
Average for Flux Tubes

$$\frac{\Delta T_i}{\Delta T_e} = 4.06 \pm 0.26$$

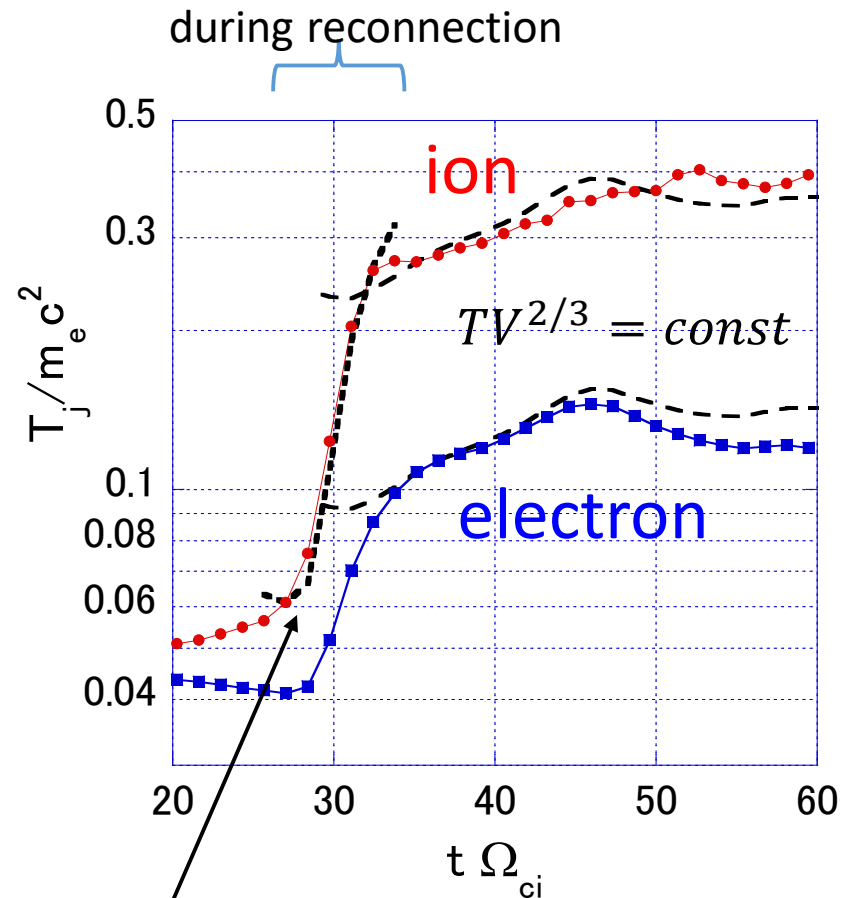
Mass dependence



Thermodynamics of Reconnection



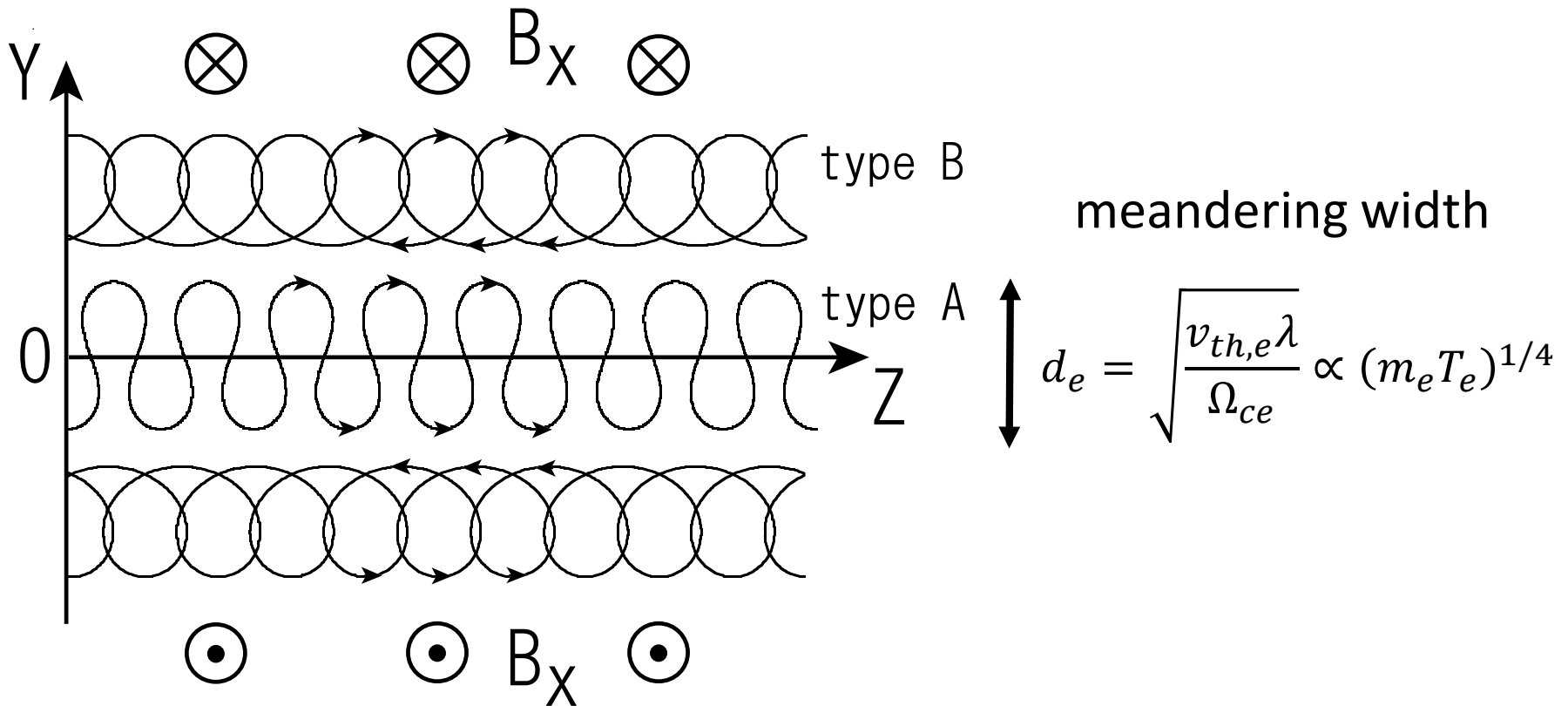
(V : Volume of Flux Tube)



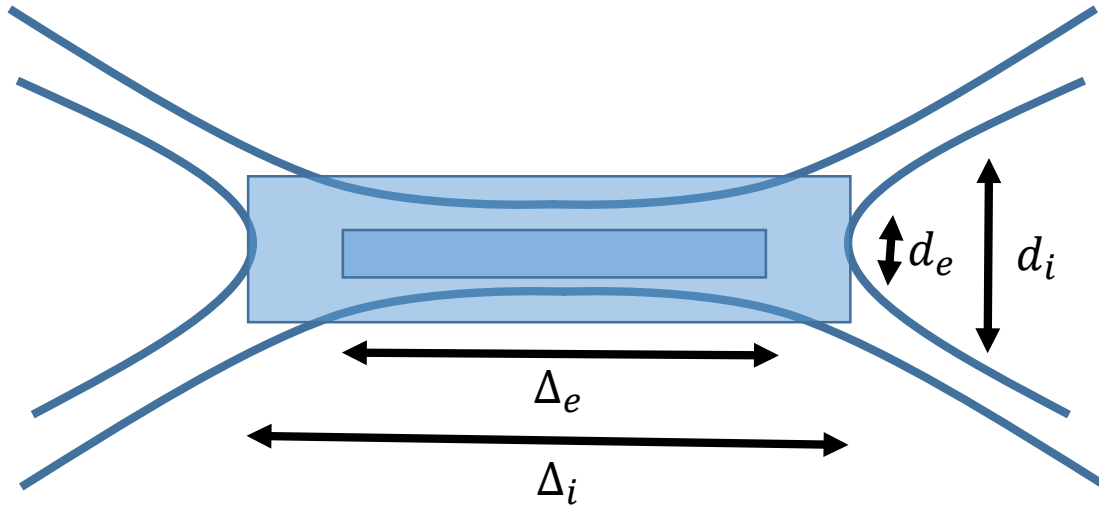
Heating during Reconnection

$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i}{m_e} \right)^{1/4}$$

meandering motion in diffusion region

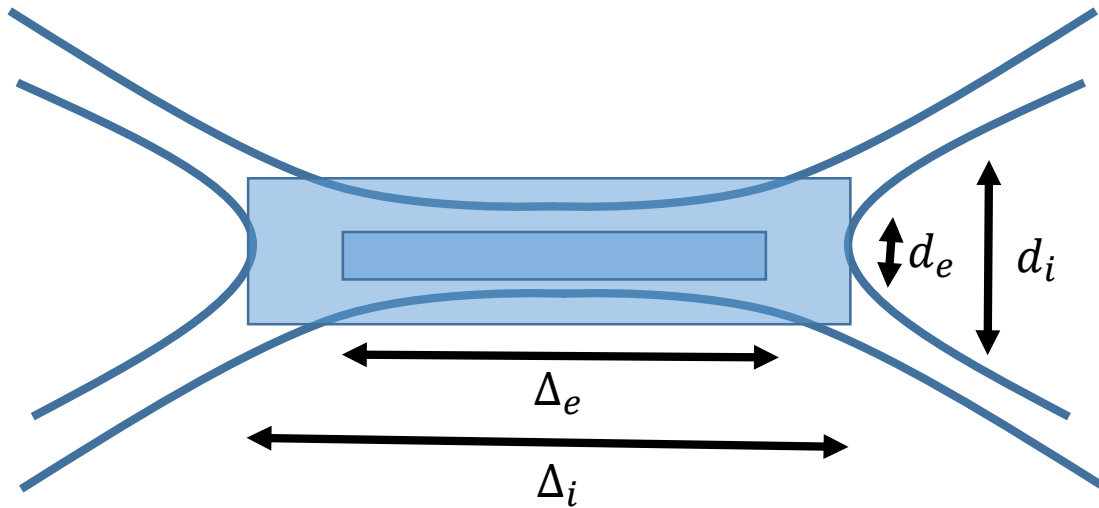


Joule heating model (I)



$$\frac{\Delta T_i}{\Delta T_e} = \frac{\text{Ion Heating}}{\text{Electron Heating}} = \frac{E \cdot J_i \Delta_i d_i}{E \cdot J_e \Delta_e d_e} = \frac{J_i \Delta_i}{J_e \Delta_e} \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

Joule heating model (II)

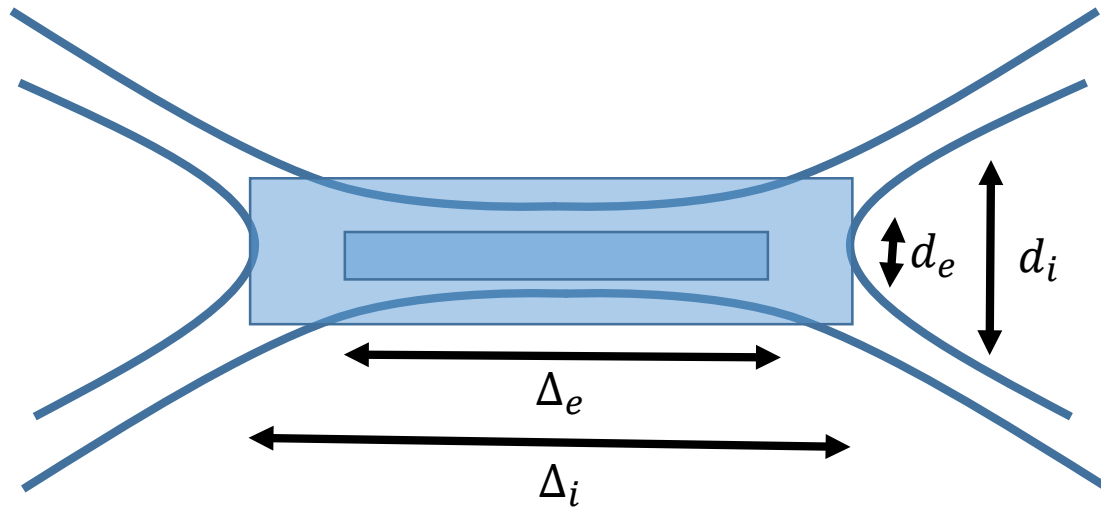


$$\frac{J_i}{J_e} = \frac{\sigma_i E}{\sigma_e E} = \frac{\frac{ne^2}{m_i} \frac{\Delta_i}{v_{ix}}}{\frac{ne^2}{m_e} \frac{\Delta_e}{v_{ex}}},$$

(e.g., Coppi, Laval & Pellat, PRL 1966; Hoh, PoF 1996)

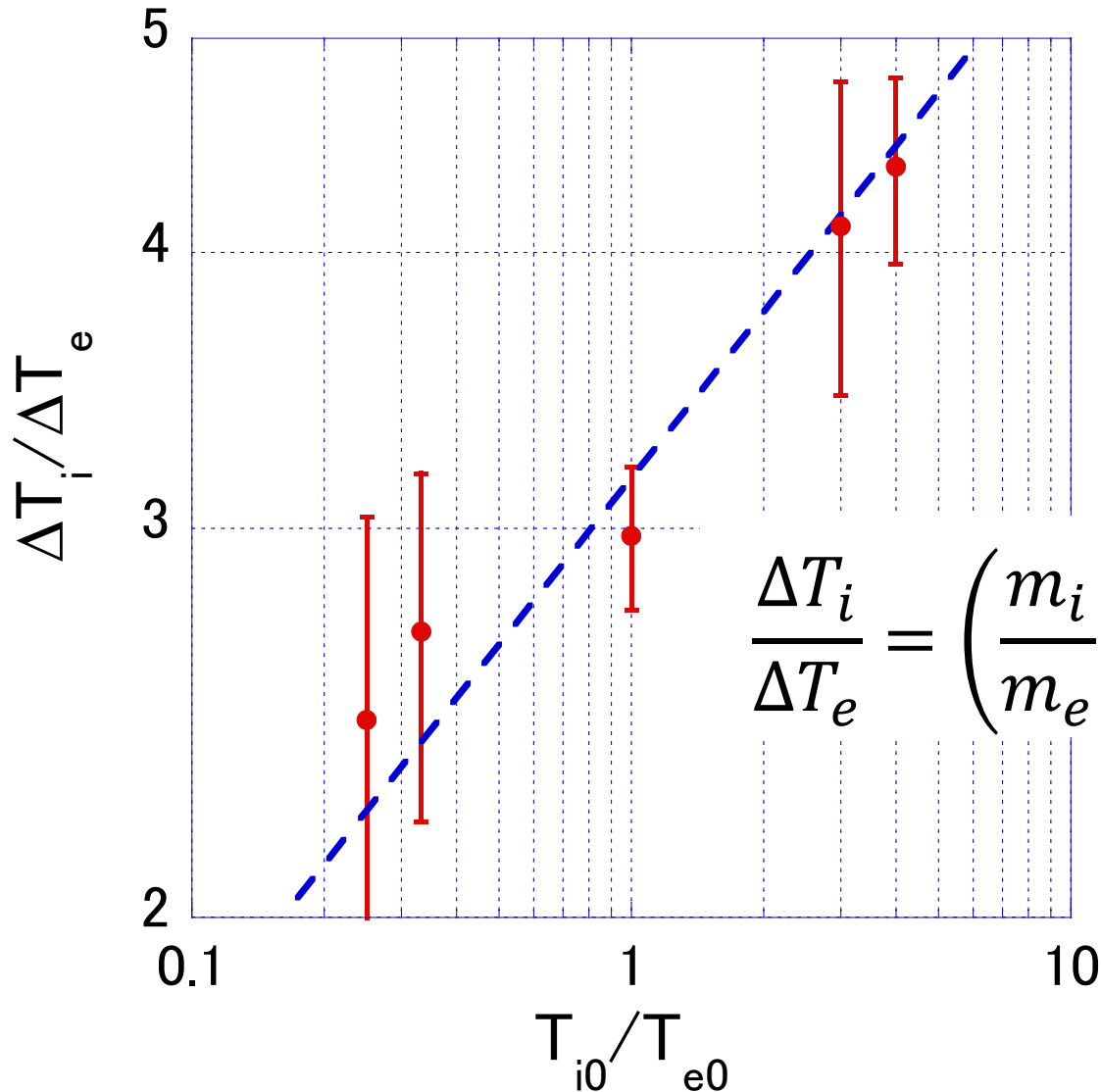
$$\frac{\Delta_i}{\Delta_e} = \left(\frac{v_{ix} \Omega_{ce}}{\Omega_i v_{ex}} \right)^{1/2} \quad \therefore \frac{J_i \Delta_i}{J_e \Delta_e} = 1$$

Joule heating model (III)



$$\frac{\Delta T_i}{\Delta T_e} = \frac{\text{Ion Heating}}{\text{Electron Heating}} = \frac{E \cdot J_i \Delta_i d_i}{E \cdot J_e \Delta_e d_e} = \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

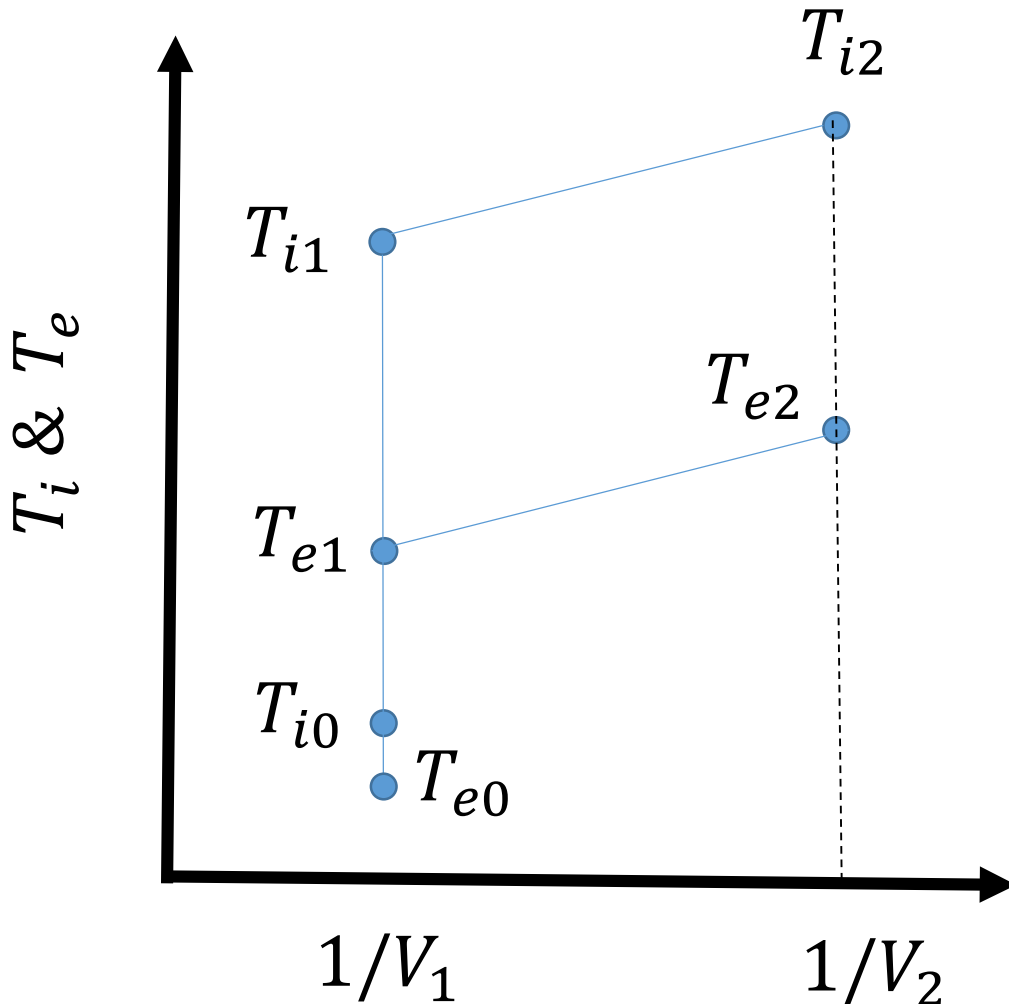
Initial temperature dependence



$$\frac{m_i}{m_e} = 100$$

$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

Thermodynamics of Reconnection



$$\frac{T_{i1} - T_{i0}}{T_{e1} - T_{e0}} = \left(\frac{m_i T_{i0}}{m_e T_{e0}}\right)^{1/4}$$

$$\frac{T_{e2}}{T_{e1}} = \left(\frac{V_1}{V_2}\right)^{\gamma-1} \quad \gamma = \frac{5}{3}$$

$$\frac{T_{i2}}{T_{i1}} = \left(\frac{V_1}{V_2}\right)^{\gamma-1}$$

Summary (Plasma Heating)

□ Energy Partition of Ion & Electron during Magnetic Reconnection

□ Two distinct heating stages:

□ Effective Ohmic heating

$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

□ Adiabatic Compression

$$\frac{D}{Dt} (TV^{\gamma-1}) = 0$$