Dark energy in the Universe



Yasushi Suto Department of Physics, The University of Tokyo Summer School of KEK theory group @ Hakuba July 30 and 31, 2007 1

Plan of the talk

- I importance of darkness
- I cosmic expansion and dark energy
- III cosmic acceleration and SN Ia Hubble diagram
- IV current constraints on dark energy
- V future dark energy projects

I importance of darkness



Blue sky at Bologna on June 23, 2007

Without dark nights, one could have never imagined ...

what really dominates our world

a planet with six Suns





no "night" except the total eclipse due to another planet every 2050 years

People realized the true world for the first time through the darkness full of "stars" Darkness is the key to understanding our world better

- Beyond the edge of our current horizon (= "darkness", "dark night")
- philosophy, astronomy, and therefore physics started from thinking in the dark
 Should still apply now
 - Another element: dark matter, dark energy
 - Another Earth: extrasolar planet
 - Another world: Multiverse
 - Another life: extra-terrestrial intelligence

Ancient particle physics





クォーク

10⁻¹⁶

cm

(いずもりよう:須藤靖「ものの大きさ」図1.1より)

SDSS (Sloan Digital Sky Survey) Apache Point Observatory @New Mexico, US



NHK education "Science Zero" broadcast on June 11, 2003

Progress of our eyes to the universe





Ground 4m telesope + CCD = 100 × photographic plate

ST Scl OPO January 15, 1996 R. Williams and the HDF Team (ST Scl) and NAS

HST WFPC

HST(2.4m)+CCD =1000 × ground telescopes

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Distant universe observed by Subaru telescope





http://hubblesite.org/newscenter/newsdesk/archive/releases/2006/23/ Light emitted from quasar bends around intervening galaxy cluster,

QSO at 10 billion light years away



Gravitational lens SDSS J1004+4112 : general relativistic mirage Galaxy cluster at 6 billion light years away

producing lensed images*

Mirage from the universe 10 billion years ago (SDSS J1004+4112)

Discovered by N.Inada and M.Oguri at Univ. of Tokyo in 2003 from SDSS images and then confirmed by Subaru and Keck Inada et al. Nature 426(2003)810

Back to the universe 10 billion years ago



HST photo release on May 23, 2006 http://hubblesite.org/newscenter/newsdesk/archive/releases/2006/23/

Galaxy Cluster SDSS J1004+4112 HST ACS/WFC

Gravitational lens SDSS J1004+4112

Lensed Galaxy



10"

Gravitationally Lensed Quasar in Galaxy Cluster SDSS J1004+4112

HST ACS/WFC

Looking toward the edge of the universe has revealed the presence of dark matter

NASA, ESA, K. Sharon (Tel Aviv University) and E. Ofek (Caltech)

STScI-PRC05-23

Dark energy in the universe



Why important ?

New physics

- major but unknown component of the universe ?
- Breakdown of general relativity at cosmological scales?

Astronomy is the key



Steven Weinberg "Right now, not only for cosmology but for elementary particle theory this is the bone in the throat"

Edward Witten *"Would be number one on my list of things to figure out"*

Frank Wilczek *"Maybe the most fundamentally ysterious thing in basic science"*

Why observable ?

Objects are usually identified only through differential observations

- Visible matter: contrast between dark and bright regions
- Dark matter: spatial inhomogeneities dynamically and gravitationally traced by visible stars, galaxies and quasars

Dark energy, if exists in a completely homogeneous manner, requires an absolute measurement for detection !?

Time variation (cosmic acceleration, structure growth): differential observation in a time, not spatial, domain

Signatures of dark energy

- cosmic acceleration
- geometry of the universeevolution of structure

Probes

- Supernova Hubble diagram
- Cosmic Microwave Background
- Gravitational lensing
- Baryon Acoustic Oscillation





I cosmic expansion and dark energy Dark Energy



Expanding the expanding universe Expand the "radius" of the universe $R(t) = R(t_0) + \frac{dR}{dt} \bigg|_{t_0} (t - t_0) + \frac{1}{2} \frac{d^2 R}{dt^2} \bigg|_{t_0} (t - t_0)^2 + \cdots$ current size:

 $R(t_0) \Leftrightarrow$ no physical meaning: $a(t) = R(t)/R(t_0)$ current expansion rate: the Hubble constant

 $H_0 \equiv \frac{dR / dt}{R}$ \Leftrightarrow unpredictable: simply due to the initial condition (can be either negative or positive)

current acceleration rate: the deceleration parameter

 $q_0 \equiv -\frac{R d^2 R / dt^2}{(dR / dt)^2} \Leftrightarrow \text{ related to the cosmic ended of the cosmic$ ⇔ related to the cosmic energy 24

The Friedmann equations

Alexander Friedmann (1888 - 1925)

the Einstein equations:

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$

 $\frac{\ddot{a}(t)}{\sigma(t)} = -\frac{4\pi G}{2} \left(\rho(t) + 3p(t)\right)$

 \Rightarrow the Friedmann equations

energy conservation equation

$$\left(\frac{\dot{a}(t)}{a(t)}\right)^{2} = \frac{8\pi G}{3}\rho(t) - \frac{K}{a^{2}(t)} + \frac{\Lambda}{3}$$
 Additional terms from on of motion

equation of motion

relativity

from



Cosmological parameters defined energy conservation $\left(\frac{\dot{a}(t)}{a(t)}\right)^2 \equiv H^2(t) = \frac{8\pi G}{3}\rho_m(t) + \frac{8\pi G}{3}\rho_r(t) - \frac{K}{a^2(t)} + \frac{\Lambda}{3}$ $\Rightarrow 1 = \Omega_m(t) + \Omega_r(t) + \Omega_K(t) + \Omega_\Lambda(t)$ density parameters matter and Λ $\Omega_m(t) \equiv \frac{8\pi G\rho_m(t)}{3H^2(t)} \quad \Omega_\Lambda(t) \equiv \frac{\Lambda}{3H^2(t)} \quad \begin{array}{c} \Pi_{\Lambda}(t) \equiv \frac{\Lambda}{3H^2(t)} \\ \text{energies in units} \end{array}$ of kinetic energy deceleration parameter $\frac{\ddot{a}(t)}{a(t)} = -\frac{4\pi G}{3} \left(\rho(t) + 3p(t) - \frac{\Lambda}{4\pi G}\right)$ $\Rightarrow q(t) = \frac{\Omega_m(t)}{2} - \Omega_{\Lambda}(t) \quad (> \text{ positive if } \Lambda = 0)$ 26

■ Simplest universe: $\Omega_m = 1, \Omega_K = \Omega_\Lambda = \Omega_r = 0$ ■ total energy = 0 (⇔ flat space): $\Omega_K = 0$

• no cosmological constant: $\Omega_{\Lambda} = 0$

■ radiation negligible: $Ω_r \doteq 0$

Deceleration parameter

$$- q = \Omega_m/2 - \Omega_{\Lambda} = 0.5$$

• q>0 should be robust as long as $\Omega_{\Lambda}=0$

 gravity is always attractive, and thus decelerates the motion. This is why the <u>deceleration</u> parameter was introduced.

Observed values

fairly complicated



- total energy \rightleftharpoons 0 (⇔ flat space, $Ω_{K} \rightleftharpoons$ 0) seems OK
- radiation negligible: $\Omega_r \doteq 0$
- something like cosmological constant (?) dominates the current universe: $\Omega_{\Lambda} = 3/4$
- $\Omega_{\rm m} \doteq 1/4$, more than 80% of the matter is dark $(\Omega_{\rm DM} \doteq 0.2, \Omega_{\rm baryon} \doteq 0.04)$

negative deceleration parameter !

- $q = \Omega_m/2 \Omega_\Lambda = -0.6 < 0$
- currently accelerating (repulsive force?)
- should have defined the acceleration parameter.

From cosmological constant to dark energy

- 1916: general relativity
- 1917: Einstein's static universe
- After 1980's: vacuum energy density

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} + \Lambda g_{\mu\nu} = 8\pi G T_{\mu\nu}$$
Cosmological constant
(geometrical quantity)

$$R_{\mu\nu} - \frac{1}{2} R g_{\mu\nu} = 8\pi G \left(T_{\mu\nu} - \frac{\Lambda}{8\pi G} g_{\mu\nu} \right)$$
Dark energy
(matter field)

Natural value: the Planck units

$$\Lambda = \frac{c^5}{\hbar G} \approx 5.2 \times 10^{93} \text{ g/cm}^3 \iff \Omega_{\Lambda} \equiv \frac{\Lambda}{3H_0^2} \approx 10^{121}$$
Observed value: $\Omega_{\Lambda} \approx 0.7$
The worst discrepancy in the history of physics ! 29

Dark energy and the equation of state of the universe Parameterized equation of state (pressure) = w x (density) ■ w=0: dark matter, ■ w=1/3: radiation w=-1: cosmological constant Poisson eq. in GR : $\Delta \phi = 4 \pi G(\rho + 3p) = 4 \pi G \rho (1 + 3w)$ $w < -1/3 \Rightarrow$ repulsion force Negative pressure: dark energy More generally w may change with time 30 w=-1 or not: that is the question conventional parameterization (no physics): $w(a) = w_0 + w_a(1-a)$ where a = 1/(1+z)cosmological constant ($w_0 = -1 & w_a = 0$)??? $w_a = 0$ or ≠0 ??? $w_a = 0$ or ≠0 ??? $w_0 = -1$ or ≠-1 ???

physical models desperately needed

 My colleagues told me that DGP (Dvali-Gabadadze-Porrati) model is approximated by

 \Rightarrow $w_0 = -0.78, w_a = 0.32$ for $\Omega_m = 0.27$ although I cannot even pronounce their names...31

 $w(a) = -\frac{1}{1 + \Omega_m(a)}$ where $\Omega_m(a) = \frac{\Omega_m}{a^3} \left(\frac{H_0}{H(a)}\right)$

Time-dependent w model • if $p = w(a) \rho$, $\rho(a) = \rho_0 \exp\left[\int_a^1 3[1+w(x)]\frac{dx}{x}\right] \equiv \rho_0 f(a)$ $H^{2}(a) = H_{0}^{2} \left[\frac{\Omega_{m}}{a^{3}} - \frac{\Omega_{K}}{a^{2}} + (1 - \Omega_{m} + \Omega_{K}) f(a) \right]$ • for $w(a) = w_0 + w_a(1-a)$

$$f(a) = a^{-3(1+w_0+w_a)} e^{-3w_a(1-a)}$$

III cosmic acceleration and SN Ia Hubble diagram

http://hubblesite.org/newscenter/archive/releases/nebula/2005/37/image/b/

Galaxy spectrum







Galaxy spectrum $= \Sigma$ (member star spectrum) Redshifted due to the cosmic expansion Recession velocity is proportional to the distance of the galaxy

(courtesy of K.Yahata)

Redshift and recession velocity of galaxies



(E.Hubble; The realm of nebulae)

Hubble's law (1929)



Estimated distance

H₀=530 instead of 70km/s/Mpc // (due to errors in estimated distance)
Type la Supernova

Progenitor: white dwarf + red giant a final stage of binary star systems



white dwarf increases its mass via accretion from the red giant

- Maximum mass of white dwarf (pressure due to the electron degeneracy > gravity)
- the Chandrasekhar mass (M_{CH}≒1.4M_{sun})
- cannot support gravity and explodes if M>M_{CH}

Light-curve of Type Ia Supernova

 peak luminosities of all SNe Ia with known distance agree within 10 percent
 discover and monitor SNe Ia for standard candles (distance indicator)



Hubble Space Telescope images



Supernova Cosmology Project: Strategy



http://www-supernova.lbl.gov/

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Supernova Cosmology Project: analysis



Detection method

- deep images of regions on the sky
- do this again one month later
- compare two sets of images, looking for new "stars" superimposed on galaxies
- Spectroscopic follow-up
 - Several types of supernovae
 - SNe la have characteristic spectra

http://www-supernova.lbl.gov/

Multicolor light-curve fitting of SNe Ia



- Brighter SNe Ia ⇒ slower decline of the peak luminosity
- Empirical scaling relation
 between the peak
 luminosity and the shape
 of the light-curve
- More accurate distance estimate after the correction using the empirical scaling

http://www-supernova.lbl.gov/ (Perlmutter 2004, Physics Today, April, p.53)

Standard candle: Type Ia Supernova

observed flux: F

absolute

SN2001cw (z=0.93)



Distance: D

dark energy parameter can be read off from the comparison between the model and the observation

SN Ia

observational estimate

Accelerating universe from SN Ia data

Dimmer (more distant)



SN Legacy Survey (Astier et al. 2006)

Constraints on $\,\Omega_{\,\rm m}$ and $\,\Omega_{\,\Lambda}$ from SN Ia



acceleration of the universe $4\pi G$ $(\rho + 3p) + \frac{\Lambda}{2}$ at the present epoch $=H_0^2 \Omega_{\Lambda}$ accelerating universe if $\Omega_{\Lambda} > \Omega_{m}/2$

IV current constraints on dark energy



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Wilkinson Microwave Anisotropy Probe (WMAP) Three Year Results: Implications for Cosmology

D.N.Spergel et al. ApJS 170(2007)377





CMB acoustic oscillation



NASA/WMAP Science Team

Photon acoustic oscillation

Photon fluid behavior inside cosmic horizon

- Oscillation due to pressure
- Θ_0 : monopole component of $\delta T/T$

$$\frac{d^2 \widetilde{\Theta}_0(k,\eta)}{d\eta^2} + \frac{1}{a} \frac{da}{d\eta} \left(1 - 3c_s^2\right) \frac{d \widetilde{\Theta}_0(k,\eta)}{d\eta} + k^2 c_s^2 \widetilde{\Theta}_0(k,\eta) \approx 0$$

η: conformal time (dt=adη)

c_s(η): sound velocity

• For adiabatic density fluctuations $\tilde{\Theta}_0(k,\eta) \approx A(k) \cos[kr_s(\eta)]$

e.g., Kodama & Sasaki (1987) 48

Sound horizon scale

comoving distance that the sound wave propagates before cosmic time t

$$\mathbf{r}_{s}(t) = \int_{0}^{t} \frac{\mathbf{c}_{s}(t)}{a(t)} dt \quad (t < t_{dec})$$

where

$$c_{s}^{2} = \frac{\partial p}{\partial \rho} = \frac{1}{3} \frac{\partial p}{\partial (\rho_{\gamma} + \rho_{b})} = \frac{1}{3(1+R)}, \quad R \equiv \frac{3\rho_{b}}{4\rho_{\gamma}} = \frac{3\Omega_{b}}{4\Omega_{\gamma}}a$$
$$\frac{da}{dt} = \frac{H_{0}}{a} \sqrt{a\Omega_{m} + \Omega_{rad}}$$
$$\Rightarrow r_{s}(t_{dec}) = \frac{4\sqrt{\Omega_{\gamma}}}{3H_{0}\sqrt{\Omega_{b}\Omega_{m}}} \ln\left(\frac{\sqrt{a + \Omega_{rad}/\Omega_{m}} + \sqrt{a + 4\Omega_{rad}/3\Omega_{b}}}{\sqrt{\Omega_{rad}/\Omega_{m}} + \sqrt{4\Omega_{rad}/3\Omega_{b}}}\right)$$
$$\approx 147(0.13/\Omega_{m}h^{2})^{0.25}(0.024/\Omega_{b}h^{2})^{0.08} \text{ Mpc}$$

Temperature fluctuation angular spectrum



Weighing the universe



dark matter

dark energy

ordinary matter makes up merely 4 percent of the entire mass of the universe

 galaxies and clusters are surrounded by invisible mass an order-of-magnitude more massive than their visible part
 unknown elementary particles?

universe is dominated by even more exotic component !
homogeneously fills the universe (unclustered)
repulsive force (negative pressure; equation of state:P=- \(\rho\))
Einstein's cosmological constant ?

WMAP 3yrs: cosmological parameters

+ flat powerlaw \Lambda CDM

 $\Omega_m h^2 = 0.127^{+0.007}_{-0.013}$

 $\Omega_b h^2 = 0.0223^{+0.0007}_{-0.0009}$

 $h = 0.73^{+0.03}_{-0.03}$

 $n_s = 0.951^{+0.015}_{-0.019}$

 $\tau = 0.09^{+0.03}_{-0.03}$

 $\sigma_8 = 0.74^{+0.05}_{-0.06}$

+flat+SNLS

 $w = -0.97^{+0.07}_{-0.09}$

• + SNLS + LSS $w = -1.06^{+0.13}_{-0.08}$

+w=-1+SNLS

 $\Omega_k = -0.015^{+0.020}_{-0.016}$

+others

 $\Sigma m_{v} < 0.68 \mathrm{eV}$

WMAP: 1st year vs. 3 years

Table 2: Power Law Λ CDM Model Parameters and 68% Confidence Intervals. The Three Year fits in this Table assume no SZ contribution, $A_{SZ} = 0$, to allow direct comparison with the First Year results. Fits that include SZ marginalization are given in Table 5 (first column) and represent our best estimate of these parameters.

Parameter	First Year	WMAPext	Three Year	First Year	WMAPext	Three Year
	Mean	Mean	Mean	ML	ML	ML
$100\Omega_b h^2$	$2.38^{+0.13}_{-0.12}$	$2.32^{+0.12}_{-0.11}$	2.23 ± 0.08	2.30	2.21	2.22
$\Omega_m h^2$	$0.144^{+0.016}_{-0.016}$	$0.134^{+0.006}_{-0.006}$	0.126 ± 0.009	0.145	0.138	0.128
H_0	72^{+5}_{-5}	73^{+3}_{-3}	74^{+3}_{-3}	68	71	73
au	$0.17^{+0.08}_{-0.07}$	$0.15_{-0.07}^{+0.07}$	0.093 ± 0.029	0.10	0.10	0.092
n_s	$0.99^{+0.04}_{-0.04}$	$0.98^{+0.03}_{-0.03}$	0.961 ± 0.017	0.97	0.96	0.958
Ω_m	$0.29^{+0.07}_{-0.07}$	$0.25^{+0.03}_{-0.03}$	0.234 ± 0.035	0.32	0.27	0.24
σ_8	$0.92^{+0.1}_{-0.1}$	$0.84^{+0.06}_{-0.06}$	0.76 ± 0.05	0.88	0.82	0.77

a factor of 2 reduction of τ
 n_s is consistent with unity
 σ₈ and Ω_m become smaller

WMAP \CDM best-fit parameters

	WMAP		WMAP		WMAP+ACBAR		WMAP +			
	Only		+CBI+VSA		+BOOMERanG		2dFGRS			
Parameter										
$100\Omega_b h^2$		$2.233^{+0.0}_{-0.0}$	72 91	$2.203^{+0.}_{-0.}$	$072 \\ 090$	$2.228^{+0.066}_{-0.082}$		$2.223^{+0.066}_{-0.083}$		
$\Omega_m h^2$		$0.1268^{+0.0}_{-0.0}$	$\frac{073}{128}$	$0.1238^{+0.}_{-0.}$	$0066 \\ 0118$	$0.1271_{-0.0128}^{+0.0070}$		0.1	$0.1262^{+0.0050}_{-0.0103}$	
h		$0.734^{+0.0}_{-0.0}$	28 38	$0.738^{+0.0}_{-0.0}$	028 037	$0.733\substack{+0.030\\-0.038}$		0	$0.732_{-0.025}^{+0.018}$	
A		$0.801^{+0.0}_{-0.0}$	$\frac{43}{54}$	$0.798^{+0.0}_{-0.0}$	$047 \\ 057$	$0.801\substack{+0.048\\-0.056}$		$0.799_{-0.051}^{+0.042}$		
au		$0.088^{+0.0}_{-0.0}$	28 34	$0.084^{+0.0}_{-0.0}$	$031 \\ 038$	$0.084\substack{+0.027\\-0.034}$		$0.083\substack{+0.027\\-0.031}$		
$n_s = 0.951^{+0.00}_{-0.00}$		$0.951^{+0.0}_{-0.0}$	$15 \\ 19$	$0.945\substack{+0.015\\-0.019}$		$0.949\substack{+0.015\\-0.019}$		$0.948^{+0.014}_{-0.018}$		
σ_8	$\sigma_8 = 0.744^{+0.09}_{-0.00}$		50 60	$0.722_{-0.056}^{+0.044}$		$0.742^{+0.045}_{-0.057}$		$0.737_{-0.045}^{+0.033}$		
$\Omega_m = 0.238^{+0}_{-0}$		$0.238^{+0.0}_{-0.0}$	27 45	$0.229_{-0.042}^{+0.026}$		$0.239\substack{+0.025\\-0.046}$		0	$0.236^{+0.016}_{-0.029}$	
	1	WMAP+	V	VMAP+	W	MAP+	WMAP -	╞	WMAP+	
		SDSS		LRG	5	SNLS	SN Gold	l	CFHTLS	
Parameter										
$100\Omega_b h^2$	2	$2.233^{+0.062}_{-0.086}$	2.3	$242^{+0.062}_{-0.084}$	2.2	$33\substack{+0.069\\-0.088}$	$2.227^{+0.06}_{-0.08}$	65 32	$2.247^{+0.064}_{-0.082}$	
$\Omega_m h^2$	0.	$1329_{-0.0109}^{+0.0057}$	0.1	$337^{+0.0047}_{-0.0098}$	0.12	$95\substack{+0.0055\\-0.0106}$	$0.1349^{+0.00}_{-0.01}$)54 106	$0.1410^{+0.0042}_{-0.0094}$	
h	0	$0.709^{+0.024}_{-0.032}$	0.'	$709^{+0.016}_{-0.023}$	0.7	$23^{+0.021}_{-0.030}$	$0.701^{+0.02}_{-0.02}$	20 26	$0.686\substack{+0.017\\-0.024}$	
A	0	$0.813^{+0.042}_{-0.052}$	0.8	$816^{+0.042}_{-0.049}$	0.8	$08^{+0.044}_{-0.051}$	$0.827^{+0.04}_{-0.05}$	15 53	$0.852^{+0.036}_{-0.047}$	
au	0	$0.079^{+0.029}_{-0.032}$	0.0	$082^{+0.028}_{-0.033}$	0.0	$85^{+0.028}_{-0.032}$	$0.079^{+0.02}_{-0.03}$	28 34	$0.088^{+0.021}_{-0.031}$	
n_s	0	$0.948^{+0.015}_{-0.018}$	0.	$951^{+0.014}_{-0.018}$	0.9	$50^{+0.015}_{-0.019}$	$0.946^{+0.01}_{-0.01}$.5 .9	$0.950^{+0.015}_{-0.019}$	
σ_8	0	$0.772^{+0.036}_{-0.048}$	0.'	$781^{+0.032}_{-0.045}$	0.7	$58^{+0.038}_{-0.052}$	$0.784^{+0.03}_{-0.04}$	35 19	$0.826^{+0.023}_{-0.035}$	
Ω_m	0	$0.266^{+0.025}_{-0.040}$	0.5	$267^{+0.017}_{-0.029}$	0.2	$49_{-0.034}^{+0.023}$	$0.276^{+0.02}_{-0.03}$	22 36	$0.301\substack{+0.018\\-0.031}$	

Constraints on w from WMAP3yr + other data



Constraints on w in flat universes

Data Set	with perturbations	no perturbations
WMAP + SDSS	$-0.75_{-0.16}^{+0.18}$	$-0.69^{+0.19}_{-0.18}$
WMAP + 2dFGRS	$-0.914^{+0.193}_{-0.099}$	$-0.877^{+0.094}_{-0.110}$
WMAP + SNGold	$-0.944^{+0.076}_{-0.094}$	$-0.940^{+0.071}_{-0.092}$
WMAP + SNLS	$-0.966^{+0.070}_{-0.090}$	$-0.984^{+0.066}_{-0.085}$
CMB+LSS+SN	$-0.926^{+0.051}_{-0.075}$	$-0.915^{+0.049}_{-0.075}$

SN legacy survey 1st year (Astier et al. 2006) + SDSS LRG BAO (Eisenstein et al. 2005) \downarrow w=-1.023±0.090(sys.)±0.054(stat.)

Constraints on w in non-flat universes



Constraints on the spatial curvature and cosmological constant

Data Set	Ω_K	Ω_{Λ}
WMAP + $h = 0.72 \pm 0.08$	$-0.003^{+0.013}_{-0.017}$	$0.758^{+0.035}_{-0.058}$
WMAP + SDSS	$-0.037^{+0.021}_{-0.015}$	$0.650^{+0.055}_{-0.048}$
WMAP $+ 2dFGRS$	$-0.0057^{+0.0061}_{-0.0088}$	$0.739^{+0.026}_{-0.029}$
WMAP + SDSS LRG	$-0.010^{+0.014}_{-0.012}$	$0.728^{+0.020}_{-0.027}$
WMAP + SNLS	$-0.015^{+0.020}_{-0.016}$	$0.719_{-0.029}^{+0.021}$
WMAP + SNGold	$-0.017^{+0.022}_{-0.017}$	$0.703^{+0.030}_{-0.038}$

Spergel et al. (2007)

The cosmological standard model: What's next?

Cosmology is now in a similar stage in its intellectual development to particle physics three decades ago when particle physicists converged on the current standard model. The standard model of particle physics fits a wide range of data, but does not answer many fundamental questions: "what is the origin of mass? why is there more than one family ?, etc.". Similarly, the standard cosmological model has many deep open questions: "what is the dark energy? what is the dark matter? what is the physical model behind inflation (or something like inflation)?" Over the past three decades, precision tests have confirmed the standard model of particle physics and searched for distinctive signatures of the natural extension of the standard model: supersymmetry. Over the coming years, improving CMB, large scale structure, lensing, and supernova data will provide even more rigorous tests of the cosmological standard model and search for new physics beyond the standard model.

Spergel et al. ApJS 148 (2003) 175

Cosmology requires new physics beyond the standard model of particle physics

The standard model of cosmology has survived another rigorous set of tests. The errors on the WMAP data at large ℓ are now three times smaller and there has been significant improvements in other cosmological measurements. Despite the overwhelming force of the data, the model continues to thrive.

The data are so constraining that there is little room for significant modifications of the basic Λ CDM model. The combination of WMAP measurements and other astronomical measurements place significant limits on the geometry of the universe, the nature of dark energy, and even neutrino properties. While allowing for a running spectral index slightly improves the fit to the WMAP data, the improvement in the fit is not significant enough to require a new parameter.

Cosmology requires new physics beyond the standard model of particle physics: dark matter, dark energy and a mechanism to generate primordial fluctuations. The WMAP data provides insights into all three of these fundamental problems:

Spergel et al. ApJS 170(2007)377

Baryon acoustic oscillation (BAO)

Photon acoustic oscillation

$\tilde{\Theta}_0(k,\eta) \approx A(k) \cos[kr_s(\eta)]$

Coupling between photons and baryons through Thomson scattering leaves an oscillatory feature in baryon density fluctuations at decoupling epoch



Eventually gravity transfers the oscillatory feature in the total matter (CDM+baryon) spectrum

Standard ruler: baryon acoustic oscillation length



Sound horizon length at recombination(=c_s×0.37Myr)
 Γ_s=147 (Ω_m h²/0.13)^{-0.25} (Ω_b h²/0.024)^{-0.08} Mpc
 Estimate the distance to the CMB last-scattering surface using the above as a standard ruler

Acoustic oscillation illustrated (1)



in the early universe, the major components of the universe, i.e., dark matter, baryons, photons, neutrinos behave as a strongly-coupled single fluid

http://cmb.as.arizona.edu/~eisenste/acousticpeak/acoustic_physics.html

Acoustic oscillation illustrated (2)



neutrinos decouple earlier and start free-streaming dark matter stays around the center due to its self-gravity baryons and photons behave as a single fluid. The central concentration induces pressure and generates an outward acoustic spherical wave

http://cmb.as.arizona.edu/~eisenste/acousticpeak/acoustic_physics.html

Acoustic oscillation illustrated (3)



After recombination (z=1000, t=0.37Myr), baryons and photons decouple. photons start freestreaming while baryons keep the acoustic features

http://cmb.as.arizona.edu/~eisenste/acousticpeak/acoustic_physics.atml

Acoustic oscillation illustrated (4)



after decoupled from photons, baryons fall into the gravitational potential due to dark matter dark matter acquires the baryon acoustic feature via their gravitational evolution

http://cmb.as.arizona.edu/~eisenste/acousticpeak/acoustic_physics.html

Evolution of density profile around a peak



http://cmb.as.arizona.edu/~eisenste/acousticpeak/acoustic_physics.html

BAO as a standard ruler $r_{s} = 147(0.13/\Omega_{m}h^{2})^{0.25}(0.024/\Omega_{b}h^{2})^{0.08} \text{ Mpc}$ • Distant measurement at different epochs • Promising methodology to observationally constrain dark energy



Picture credit: Bob Nichol

Acoustic scales and geometry of the universe



NASA/WMAP Science Team

Power spectrum of mass density fluctuations with baryon acoustic oscillation effect



Evolution of baryon acoustic oscillations



T. Nishimichi

Acoustic oscillations detected



 $r_s = 147(0.13/\Omega_m h^2)^{0.25}(0.024/\Omega_b h^2)^{0.08}$ Mpc


Ω_m=0.24 best-fit WMAP model

Percival et al. (2007)

Combined constraints from SN and BAO



Recent inspirations from brane-world scenario on modified gravity

cosmic acceleration:

induced by dark energy or by extra-dimension ? matter content or law of physics ?

an example: the DGP model; gravity leaking to extra dimensions <u>"modified" Friedmann equation</u>

$$H^{2} = H_{0}^{2} \left[\Omega_{k} (1+z)^{2} + \left(\sqrt{\Omega_{M} (1+z)^{3} + \Omega_{rc}} + \sqrt{\Omega_{rc}} \right)^{2} \right] \qquad \Omega_{rc} \equiv \frac{1}{4r_{c}^{2}H_{0}^{2}}$$

"modified" Newton Potential

$$V(r) = -\frac{G_{(4)}}{r} \left[1 + \frac{2}{\pi} \left\{ -1 + \gamma + \ln\left(\frac{r}{r_c}\right) \right\} \left(\frac{r}{r_c}\right) + O(r^2) \right] : r \ll r_c \sim \frac{1}{H_0}$$

Dvali, Gabadadze & Porrati , PLB 485 (2000) 208 Deffayet, Dvali & Gabadadze, PRD 65 (2002) 044023

modified gravity vs. cosmological constant: from SDSS to WFMOS Yamamoto, Bassett, Nichol, Suto & Yahata PRD 74(2006)063525

modified Friedmann equation (spatially flat)

$$H^{2} - \frac{H^{2/n}}{r_{c}^{2-2/n}} = \frac{8\pi G}{3}\rho$$

■ n=2: DGP model, n=∞ : cosmological constant

r_c: key parameter ~1/H₀
 r<r_c: 4D space-time, r>r_c: 5D space-time
 if spatially flat (H₀r_c)^{2/n-2} = 1-Ω_m

Λ vs. the modified DGP model



ratios relative to the Λ model (spatially flat) Yamamoto et al. (2006)

Predicted apparent shifts of BAO peaks



purely linear theory, observation in Λ CDM assumed Yamamoto et al. (2006) 78

Current constraints from the SDSS LRG sample



fit to linear theory for k<0.2hMpc⁻¹ observation in ∧ CDM assumed

Yamamoto et al. (2006)

Expected constraints from future WFMOS z=1 sample



Yamamoto et al. (2006)

Current constraints on deviations from Newton's law of gravity

 $\pm \alpha \exp$

Assume the Yukawa-type deviation:

 $V(r) = -G \frac{m_1 m_2}{m_1 m_2}$



weak, if any, constraints on cosmological scales so far...

E.G. Adelberger et al. Ann.Rev.Nucl.Part.Sci. 53 (2003) 77

Empirical constraints on deviations from Newton's law of gravity via SDSS galaxy P(k)

ad-hoc and empirical approach (Shirata et al. 2005,2006)

adopt the standard Friedmann model (i.e, ACDM)
 but with <u>an additional Yukawa term</u> to gravity
 adopt the standard interpretation of CMB
 anisotropy as the initial condition for the primordial fluctuations

assume <u>scale-independent bias of SDSS</u> <u>galaxies</u>

Yukawa-type additional gravitational potential

$$V(r) = -G \int d^3r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \left[1 + \mathcal{O}\left(1 - e^{-\frac{|\mathbf{r} - \mathbf{r}'|}{2}}\right) \right]$$



small-scale: Newtonian gravity $r \ll \lambda$: $V(r) \rightarrow -G \int d^3 r' \frac{\rho(r')}{|r-r'|}$ large-scale: $G \Rightarrow G(1+\alpha)$ $r \gg \lambda$: $V(r) \rightarrow -G(1+\alpha) \int d^3 r' \frac{\rho(r')}{|r-r'|}$

stronger (weaker) gravity on large scales if $\alpha > 0$ ($\alpha < 0$), while cosmic expansion is dictated by "correct" G₈₃

Method (Shirata et al. 2005)

1) directly solve the linear perturbation equation under the modified Newtonian potential:

assu

$$\ddot{\delta}_{k} + 2H\dot{\delta}_{k} - 4\pi G\overline{\rho}\delta_{k} \left[1 + \alpha \frac{(a/k\lambda)^{2}}{1 + (a/k\lambda)^{2}}\right] = 0$$

Iming the initial conditions of
$$\delta_{k}(a_{ini}) = \delta_{k,\Lambda CDM}(a_{ini}), \quad \frac{d\delta_{k}}{da}\Big|_{a=a_{ini}} = \frac{d\delta_{k,\Lambda CDM}}{da}\Big|_{a=a_{ini}}$$

2) apply the nonlinear correction using the Peacock-Dodds formula

3) Compare the model predictions with SDSS galaxy P(k) assuming linear bias (0.01<k[h⁻¹Mpc]<0.3)

Nonlinear correction for power spectrum applying the Peacock-Dodds fit



Shirata, Shiromizu, Yoshida & Suto: Phys.Rev.D 71(2005) 064030

Comparison with SDSS galaxy P(k)



Constraints on model parameters $V(r) = -G \int d^{3}r' \frac{\rho(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} \left[1 + \alpha \left(1 - e^{-\frac{|\mathbf{r} - \mathbf{r}'|}{\lambda}} \right) \right]$



Shirata, Shiromizu, Yoshida & Suto: Phys.Rev.D 71(2005) 064030

 $\lambda = 5h^{-1}Mpc \implies -0.5 < \alpha < 0.6 \quad (3\sigma \text{ limits})$ $\lambda = 10h^{-1}Mpc \implies -0.8 < \alpha < 0.9 \quad (3\sigma \text{ limits})$

future dark energy projects





Did we make progress at all?

Egypt





Indian





Why can we conclude that this is a better picture before knowing the nature of dark matter and dark energy? 89

Towards understanding of the universe

- 1. the n-th order parameterized model of the universe
 - $\square \quad \Omega_{\Lambda'} \ \Omega_{\mathsf{m}'} \ \Omega_{\mathsf{b}'} \ \mathsf{h}, \ \sigma_8 \dots$
- 2. improve the precision/accuracy of the numbers
- 3. understand why
 - (variants of) inflation, superstring, brane...
- 4. look for something that cannot be described in the n-th order model

• w=-1 \Rightarrow w=w₀+w₀(1-a) \Rightarrow w(a) \Rightarrow w(a,r)

- Iinear bias \Rightarrow nonlinear bias \Rightarrow non-deterministic bias
- virialized spherical halo ⇒ triaxial
 ⇒ shocked+ magnetic+conductive+cosmic ray
- 5. repeat the above steps 1∼4 until you become tired (or retire) for n=1,2,3,4,5...

Can we understand the dark side of the universe in foreseeable future ?

Dark matter

maybe new results from on-going experiments in next 5-10 years, but not from astronomy

Dark energy

- unlikely to have any breakthroughs from future experiments and/or theories in high energy physics in this century
- astronomy is the key !

Dark baryons

only astronomical observations can make a scientific new contribution since high-energy physicists already know baryons too well !

Dark matter

- high-energy experiments in near future are very promising
- still room for cosmology to help understanding dark matter ?
 - density profile of dark matter halos
 - what is the "true" density profile ? core vs. cusp
 - modeling substructure statistics
 - non-spherical modeling
 - baryonic effect; star formation, feedback,,,

Dark energy

- Meaningful theoretical breakthroughs are unlikely during this century
 - \Rightarrow observational approaches are the keys !
- something really there or just virtual ?
 - right-hand-side in the Einstein equation
 - modified gravity theory
- already (too) many proposals for future observational projects
 - need more accurate modeling
 - need to control systematic effects

Probes of Dark Energy (S.Bridle)

Cosmic Shear

Evolution of dark matter perturbations Angular diameter distance Growth rate of structure

Features in matter P(k)

Standard ruler Angular diameter distance

Supernovae S

Standard candle Luminosity distance Browned and







Cluster counts Evolution of dark matter perturbations Angular diameter distance Growth rate of structure

CMB

Snapshot at ~400,000 yr, viewed from z=0 Angular diameter distance to z~1000 Growth rate of structure (from ISW)

Surveys to measure Dark Energy (S.Bridle)

2005	2010			2015
Imaging CFHTLS SUBARU	J DES, HSC	DUNE	LSST	SKA
SDSS ATLAS KIDS	VISTA Pan-STARRS		JDEM/ SNAP	
Spectroscopy FMOS				SKA
SDSS ATLAS				
Supernovae CSP ESSENCE	DES		LSST	
CFHTLS	Pan-STARRS		JDEM/ SNAP	
Clusters AMI APEX SP	T DES			
XCS SZA AMIBA AC	T			
CMB WMAP 3	WMAP 6 yr			
	Planck	Planck 4yr	•	
2005	2010			2015

 $H^{2}(z) = H^{2}_{0} \left[\Omega_{M} (1+z)^{3} + \Omega_{DF} (1+z)^{3} (1+w)\right]$ (flat) dark energy (constant w) matter Geometric tests: $r(z) = \int dz/H(z)$ Comoving distance $d_{I}(z) = (1+z) r(z)$ Standard Candles $d_{A}(z) = (1+z)^{-1} r(z)$ **Standard Rulers** $dV/dzd\Omega = r^2(z)/H(z)$ Standard Population (volume) Structure based-tests: The rate of growth of structure determined by H(z), by any modifications of gravity on large scales, and by other cosmological parameters 96

Probing Dark Energy (J.Frieman)

Probe dark energy through the expansion history:

future dark energy survey projects

DES: Dark Energy Survey (Fermi Lab+, 2011-?)

- Imaging galaxy survey
- 5000 deg²@Chile 4m telescope

HSC: Hyper Suprime-Cam (Subaru+Princeton, 2011-)

- Imaging galaxy survey 1.5deg FOV
- 2000 deg²@Subaru 8m telescope

LSST: Large Synoptic Survey Telescope (SLAC+, 2014-?)

- Imaging galaxy survey
- 20000 deg²@Chile 8.4m dedicated telescope
- WFMOS: Wide Field Multi-Objects Spectrograph (Subaru+Gemini+???, 2015-???)
 - Spectroscopic galaxy survey 1.5deg FOV
 - 4000 fibers, 20000 galaxy redshifts a night



The Dark Energy Survey (J.Frieman) Study Dark Energy using 4 complementary techniques: **Cluster counts & clustering** Weak lensing Galaxy angular clustering **SNe la distances** Two multiband surveys: 5000 deg² g, r, i, z & 40 deg² repeat (SNe) Build new 3 deg² camera Construction 2005-2009 Survey 2009-2014 (525 nights) **Response to NOAO AO** Blanco 4-meter at CTIO

Hyper Suprime-Cam project



計画研究B01(名古屋大理論):銀河分布を用いた ダークエネルギーの研究 計画研究B02(東北大理論):重カレンズ効果による暗黒物質分布 と宇宙の構造形成史の解明

公募研究:超新星探査とダークエネルギー性質解明に関する理論 および観測的研究 Ministry of Education, Special Priority Area Grant-in-Aid: 2006-2011 "Study of Dark Energy from Wide-Field Deep Survey of the Universe"

Constraining dark energy via gravitational lensing survey

- PI: Hiroshi Karoji (NAOJ)
- CCD: Satoshi Miyazaki (NAOJ)
- DAQ: Hiroaki Aihara (U.Tokyo)
- Theory groups at NAOJ, Univ. of Tokyo, Nagoya Univ. Tohoku Univ.
- Princeton Univ. will join officially

International Research Network for Dark Energy (JSPS, core-to-core program 2007-2009)



WFMOS proposal: Subaru+Gemini spectroscopic survey

- Observational constraints on dark energy
- Accurate measurement of the baryon acoustic scales in galaxy distribution
- 4000 multi-fiber spectrograph on 1.5deg FOV cameta at Subaru prime focus
 - 0.5<z<1.3: emission line galaxies</p>
 - 2×10^6 gals/2000 deg² \Rightarrow 1400 pointings(900hours)
 - 2.3<z<3.3: Lyman-break galaxies</p>
 - 6×10^5 gals/300 deg² \Rightarrow 200 pointings(800hours)
- Determine H(z) and D(z) within 1% precision
- Determine w within 3% precision and dw/dz within 25% precision

Dark energy research is good or bad for astronomy ?

Fundamentalist physics: why dark energy is bad for astronomy Simon D.M. White, astro-ph/0704.2291

Fundamentalist: high-energy experiments

- Pursuit of a single truth (LHC, WMAP)
- Huge international collaborations
- Universalist: astronomical observations
 - Multi-purpose (Hubble Space Telescope, SDSS)
 - Relatively small groups

Different culture, sense of value, and matter of taste, after all...

Two very successful but quite different projects in astronomy

HST (universalist)

An observatory Designed for general tasks Serving a diverse community Programme built through proposals Many teams of all sizes Many results unanticipated Nourishes astrophysics skills

Public support as a facility

WMAP (fundamentalist)

An experiment Designed for a specific task Serving a single, coherent community Programme set at design A single moderately large team Main results 'planned' Nourishes data-processing/ statistics skills Public impact through results

Simon D.M. White: astro-ph/0704.2291 ¹⁰³

Another successful example: SDSS (Sloan Digital Sky Survey)

- 1345 refereed papers to date
- These papers have been cited over 39,000 times
- 30 of the 200 most cited papers in astronomy since 2000 used SDSS data
- Impact in many areas we didn't anticipate:
 - White dwarfs
 - Brown dwarfs
 - Ultra-low metallicity stars
 - Galaxy-galaxy lensing
 - Supernovae
 - Epoch of reionization

(Dec. 19, 2006@NAOJ, Michael Strauss)

ADS High-Impact Papers 2006

Facility	Number of Citations	Fraction of the Total
SDSS	1843	17.4%
ESO	1365	12.9%
HST	1124	10.6%
WMAP	1121	10.6%
Keck	642	6.0%
Kamiokande	372	3.5%
Chandra	365	3.4%
ACBAR	207	2.0%
NOAO (KPNO/CTIO)	202	1.9%
Las Campanas	176	1.7%

(Dec. 19, 2006@NAOJ, Michael Strauss)



The town mouse and the country mouse (都会のねずみと 田舎のねずみ)



Town mouse ?	Country mouse ?
Large Hadron Collider	Subaru telescope
Tristan	Kamiokande
particle theorists	astronomers
dark energy cosmology	extrasolar planet
Subaru (8.2m)	HATnet (11cmx6)

So, what's next?

Precision cosmology, not yet ?

- We have to move on; determine all the cosmological parameters within 0.1% accuracy, for instance.
- For what ? Really interesting ? Can convince taxpayers ?

Beyond *precision* cosmology ?

- Stop playing with the <u>values of parameters</u>, but try to understand <u>their meaning</u>, i.e., matter context in the universe
 - Nature of dark matter and dark energy
 - First objects in the universe
 - initial conditions (physical model of inflation)...
- Revisit the cosmological observations in a more general framework
 - Equation of state of the universe
 - Validity of the cosmological principle
 - Validity of general relativity on cosmological scales

Or simply <u>beyond cosmology</u> itself !

Anthropic principle, Extrasolar planet, ...something else⁰⁷