Density profiles and clustering of dark matter halos and clusters of galaxies

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Why density profiles of dark halos ?

- Theoretical interest: what is the final state of the cosmological selfgravitating system ? • forget cosmological initial conditions? •keep initial memory somehow? Practical importance: testing cosmology and/or nature of dark matter
 - galactic rotation curve, gravitational lensing, X-ray/SZ observations of clusters

Brief history (before NFW)

- <u>1970:</u> Peebles; N-body simulation (N=300).
- <u>1977:</u> Gott; secondary infall model r -9/4.
- <u>1985:</u> Hoffman & Shaham; predict that density profile around density peaks is r^{-3(n+3)/(n+4)}.
- <u>1986:</u> Quinn, Salmon & Zurek; N-body simulations (N ~ 10000), confirmed r ^{-3(n+3)/(n+4)}.
- <u>1988:</u> Frenk, White, Davis & Efstathiou;N-body simulations (N=32³), showed that CDM model can reproduce the flat rotation curve out to 100kpc.
- <u>1990:</u> Hernquist; proposed an analytic model with a central cusp for elliptical galaxies r⁻¹(r+r_s)⁻³.
- <u>1996:</u> Navarro, Frenk & White; universal density profile for dark matter halos.



Importance of high-resolution simulations

Iow mass/force resolutions shallower potential than real artificial disruption/overmerging (especially serious for small systems)

$\varepsilon = 1 \text{kpc}$



$\varepsilon = 7.5 \text{kpc}$



Moore (2001)

central 500kpc region of a simulated halo in SCDIM

Profiles in higher-resolution simulations



Origin of the universal density profiles ?

- <u>1977:</u> Davis & Peebles; stable clustering solution of 2pt correlation function r -3(n+3)/(n+5)
- <u>1977:</u> Gott; secondary infall model r -9/4
- 1985: Hoffman & Shaham; mass profile around density peaks r^{-3(n+3)/(n+4)}
- 1997: Syer & White; dynamical friction of satellites halos
 r -3(n+3)/(n+5)

Gallery of high-resolution simulated halos

galaxies ~ 5x10¹²M_{sun}

groups ~ 5x10¹³M_{sun}

 $\frac{clusters}{\sim 3 x 10^{14} M_{sun}}$



Jing & Suto (2000)

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Weak mass-dependence of halo profiles ?

Inner slope of the profile (= 1.2 ~ 1.5) is weakly dependent on the halo mass ?



Jing & Suto (2000), but see also Fukushige & Makino (2001)

Summary of simulation and theory

Simulations

- Density profiles of dark halos are fairly universal (at least approximately), and are insensitive to the cosmological initial conditions
- Cusp rather than core in the central region

$$\rho(r) = \frac{\delta_c \rho_{crit}}{(r/r_s)^{\alpha} (1+r/r_s)^{3-\alpha}} \quad \alpha \approx 1.5$$

<u>Theoretical models</u>

- Either core or cusp is acceptable.
- Inner slope is generally expected to depend on the primordial spectrum of fluctuations.

needs observational confrontation

Rotation curves of DM dominated galaxies



dwarf spirals to giant low surface brightness galaxies indicate the central cores rather than cusps ! inconsistent with CDM simulations (Moore et al. 1999; de Blok et al. 2000; Salucci & Burkert 2000)

Gravitational lensing of CL0024+1654

HST image



Z=0.39, $L_X=5\times10^{43}$ h⁻² erg/s

reconstructed mass distribution (with 512 parameters)



Tyson, Kochanski & Dell'Antonio (1998)

Reconstructed mass profile of CL0024+1654



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Crisis of cold dark matter ?

Observations favor the presence of core rather than cusp

- Rotation curves of low-surface brightness galaxies
- Cluster mass profile from gravitational lensing
- still controversial, but ...
- Cold dark matter is really collisionless ?
 - Self-interacting dark matter

(Spergel & Steinhardt 1999)

Other hydrodynamical/radiative processes ?

- Supernova feedback
- Bar-driven core formation (Weinberg & Katz 2001)



$$(mn)\frac{\sigma}{m}\ell = 1 \quad \Rightarrow \quad \frac{\sigma}{m} = 2\text{cm}^2/\text{g}\left(\frac{10^4\rho_{\text{crit}}}{\rho_{\text{center,cl}}}\right)\left(\frac{1\text{Mpc}}{\ell}\right)$$

Collisional Dark Matter

 σ (fluid limit), steeper cusp !
 σ/m~1 cm²/g no cusp, rather forms a central core, but the resulting halos are too spherical...













 $\begin{array}{c} {\bf S1Wb} \\ \sigma^{*} = 1.0 \ {\rm cm}^{2} {\rm g}^{-1} \\ r_{\rm c} = 100 \ h^{-1} {\rm kpc} \\ 1 \ : \ 0.91 \ : \ 0.72 \end{array}$



 $\begin{array}{l} {\bf S1Wc} \\ \sigma^* = 10.0 \, {\rm cm}^2 {\rm g}^{-1} \\ r_{\rm c} = 160 \, h^{-1} {\rm kpc} \\ 1 \, : \, 0.98 \, : \, 0.89 \end{array}$



1 : 0.82 : 0.65



generally favor a steep cusp (

Constraints from the existing arc samples

Itentative application to 13 galaxy clusters with S_X>10⁻¹² erg/s/cm² and 0.1<z_L<0.4 ⇒ N_{tot, tan}=15, N_{tot, rad}=2 (Luppino et al. 1999)

Observed highfrequency of radial arcs favors the steep central cusp in massive halos as indeed suggested by CDM simulations (Molikawa & Hattori 2001)

Oguri et al. (2001)



From density profiles to clustering of dark matter: dark matter halo approach



- Accurate modeling of nonlinear clustering of dark matter
 - interpolation of linear theory and stable solution using N-body data (e.g., Hamilton et al. 1991; Peacock & Dodds 1996)
 - dark matter halo approach: pairs of particles (in a single halo + in two different halo) weighted over the halo mass function (e.g., McClelland & Silk 1977; Seljak 2000; Ma & Fry 2000)

Both results agree well when adopting the halo profiles from N-body simulations

Seljak (2000)

Clustering of luminous objects on the light-cone



2001



Predicting the clustering of dark matter on the light-cone

redshift-space distortion



average over the light-cone



Yamamoto & Suto (1998) ; Hamana, Colombi & Suto (2001) 21/32

Phenomenological model
for scale- and mass-dependent halo biasing

$$\frac{mass-dependence (Jing1998; Sheth & Tormen1999) + scale-dependence (Taruya & Suto 2000)b_{halo}(M,R,z) = b_{ST}(M,z)[1+b_{ST}(M,z)\sigma_{mass}(R,z)]^{0.15}\xi_{halo}(M,R,z) = b_{Balo}^{2}(M,R,z) \xi_{mass}(R,z)$$

$$\frac{average over the light-cone}{\xi_{halo}^{LC}(>M,r)} = \frac{\int_{z_{min}}^{z_{max}} dz \int_{M}^{\infty} dM \xi_{halo}(M,R,z) n_{ST}^{2}(M,z) \frac{dV_{c}}{dz}}{\int_{z_{min}}^{z_{max}} dz \int_{M}^{\infty} dM n_{ST}^{2}(M,z) \frac{dV_{c}}{dz}}{dz}$$

Hamana, Yoshida, Suto & Evrard; ApJ 561(2001)L143



Calibrating the halo biasing model with the Hubble volume simulation at z=0



Correlation functions of halos on the light-cone



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Summary of the current results Halo density profiles: still controversial • LSB/dwarf galaxies, CL0024-1654: a flat core N-body simulations, gravitational lensing : a cusp Needs further work from different aspects Halo clustering: a phenomenologically successful model on the light-cone • gravitational nonlinear evolution redshift-space distortion •mass-, time-, and scale-dependent bias • selection function evolution in the survey volume itself

An incomplete list of unresolved issues

Halo density profile:

- physical explanation of the central cusp
- characterizing the degree of non-sphericity
- From dark halos to visible objects:
 - halo mass -- cluster gas temperature relation
 - non-gravitational effects inside dark halos (cooling, star/galaxy formation, preheating, supernova feedback, etc.)

Clustering:

- predicting the clustering of X-ray/SZ clusters, QSOs, and massive halaxies at high-z on the light-cone.
- developing the higher-order clustering statistics using the dark halo approach

Relation between dark halos and clusters



Globally similar distribution, but their precise relation is unclear because definitions of clusters (especially at high z) are very ambiguous.

SPH simulations in LCDM: 75x75x15 h⁻³Mpc³ (Yoshikawa, Taruya, Jing & Suto 2001)

From dark halos to galaxy clusters ?

Definitely they are closely related, but the exact one-to-one correspondence is unlikely....



An example; substructure of RXJ1347-1145 (z=0.45) detected via SZ map at 150 GHz



150GHz with NOBA (Nobeyama Bolometer Array) at Nobeyama 45m telescope in March, April, 1999 and February 2000 FWHM=13"

Globally similar morphology to the X-ray image

Substructure in the South-East direction

mJy/bean

Komatsu et al. PASJ 53(2001)57

Confirmed by Chandra and BIMA observations



RXJ1347-1145

BIMA@30GHz 63"x80" beam (10.3mJy point source removed)

Carlstrom et al. (2001)

Keck spectroscopy: Cohen & Kneib (2002)
Chandra: Allen, Schmidt & Fabian (2002)
non-spherical modeling is crucial, perhaps at high z in particular.

13^h47^m36^s 32^s 28^s 24^s Right Ascension (J2000)

Triaxial model for dark halos



$$\rho(\mathbf{R}) = \frac{\delta_c \rho_{crit}}{(\mathbf{R} / \mathbf{R}_s)^{\alpha} (1 + \mathbf{R} / \mathbf{R}_s)^{3 - \alpha}}$$
$$\mathbf{R}^2(\rho) \equiv \frac{X^2}{a^2(\rho)} + \frac{Y^2}{b^2(\rho)} + \frac{Z^2}{c^2(\rho)}$$

Jing & Suto (2002) ApJ, August issue

Non-spherical description is becoming crucial in properly interpreting recent highangular resolution data of the weak/strong lensing, X-ray/SZ cluster observations

X-ray gas profiles of clusters



 isothermal gas profile in NFW potential is close to the isothermal β model
 predicted core radius is smaller than those observed



Tangential and radial arcs



Hammer et al. (1997)

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Model for halo density profile

Halo density profile

$$\rho(r) = \frac{\rho_{\rm crit} \delta_{\rm c}}{\left(r/r_{\rm s}\right)^{\alpha} \left(1 + r/r_{\rm s}\right)^{3-\alpha}}$$



$$c_{\rm vir}(M,z) = \frac{r_{\rm vir}(M,z)}{r_{\rm s}(M,z)}$$



$$c_{\rm vir}(M,z) = c_{\rm norm} \frac{2-\alpha}{1+z} \left(\frac{M_{\rm vir}}{10^{14}h^{-1}M_{\odot}}\right)^{-0.13}$$

Log-normal distribution for scatter in c_{norm}

 $\Delta(\log c_{vir})=0.18$ (Bullock et al. 2001; Jing 2000) <u>Free parameters: c_{norm} and α </u>

Expected number of arcs

<u>Number of arcs per unit solid angle</u>



Time-delay in QSO multiple images to probe the halo density profile



Time-delay is very sensitive to the inner slope, but insensitive to cosmological parameters (except H_0 !) **Steeper inner profile** larger time-delay 0.5 10-3 10-2 10-1 101 108

∆t vr

Tentative applications to 4 lens systems



Observed timedelay is consistent with predicted time-delay probability when the density profile has a steep cusp r^{-1.5}

Oguri et al. 2002

Evaluating the particle discreteness effect using dark halo approach to clustering

